

A Theoretical Model to Explain TeV Gamma-ray and X-ray Correlation in Blazars

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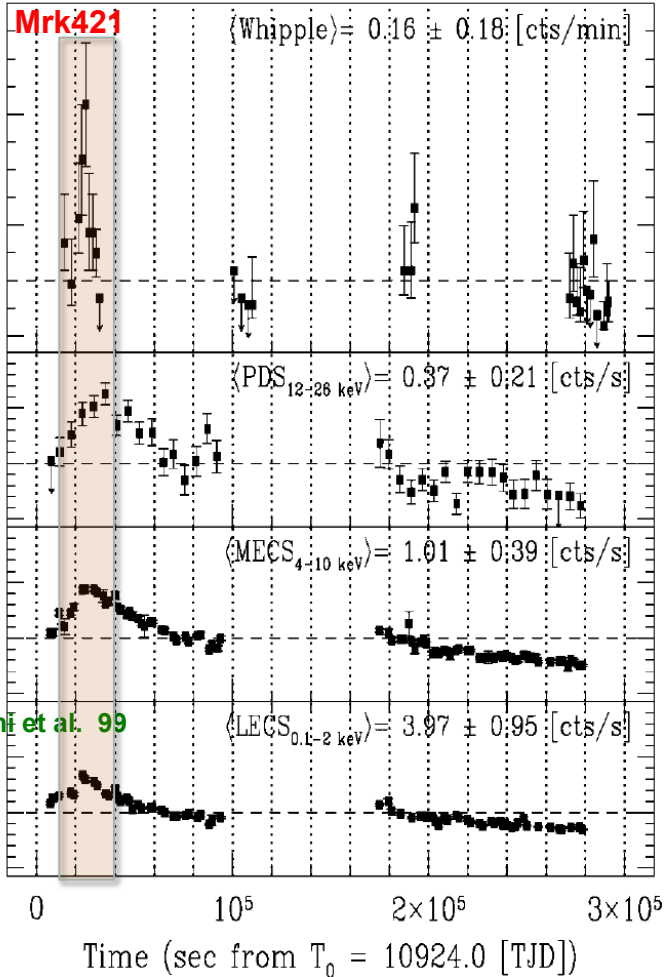
OUTLINE

- Observations: TeV and X-ray correlations and disconnections
- Theoretical model:
 - Generalities
 - Lepton model
 - Hadronic model
- Application to Mrk 421
- Conclusions

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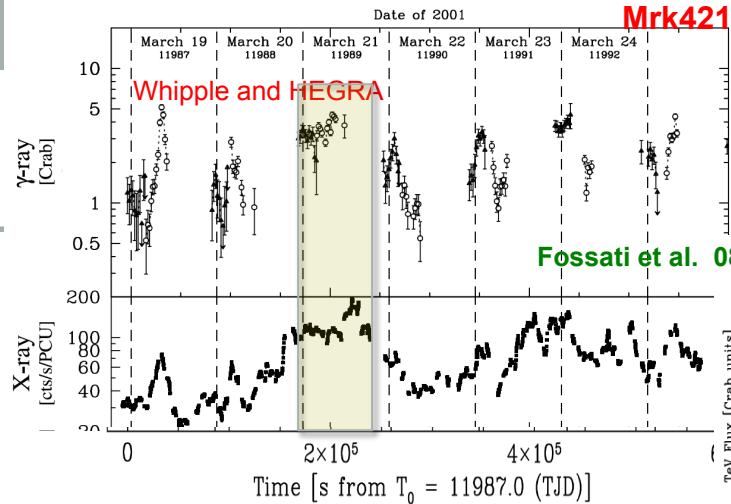
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➤ Correlations

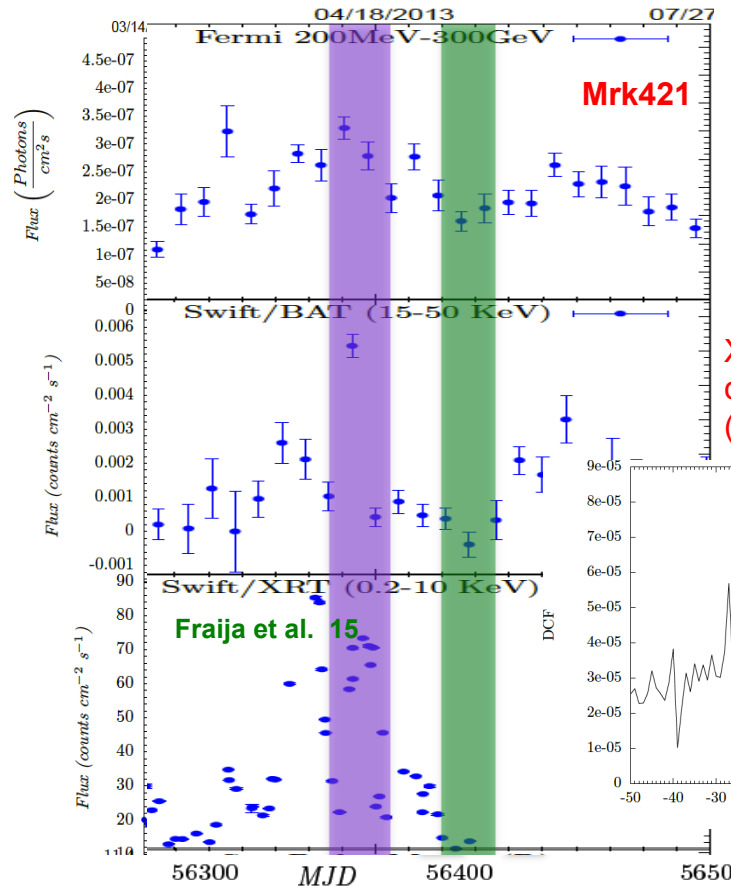
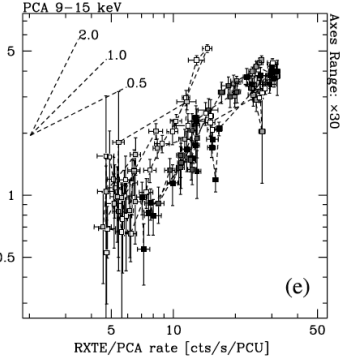


Maraschi et al. 99

X-ray and TeV are well correlated on timescale of hours (flaring activity April/98).



X-ray and TeV are highly correlated. (flaring activity 2001).



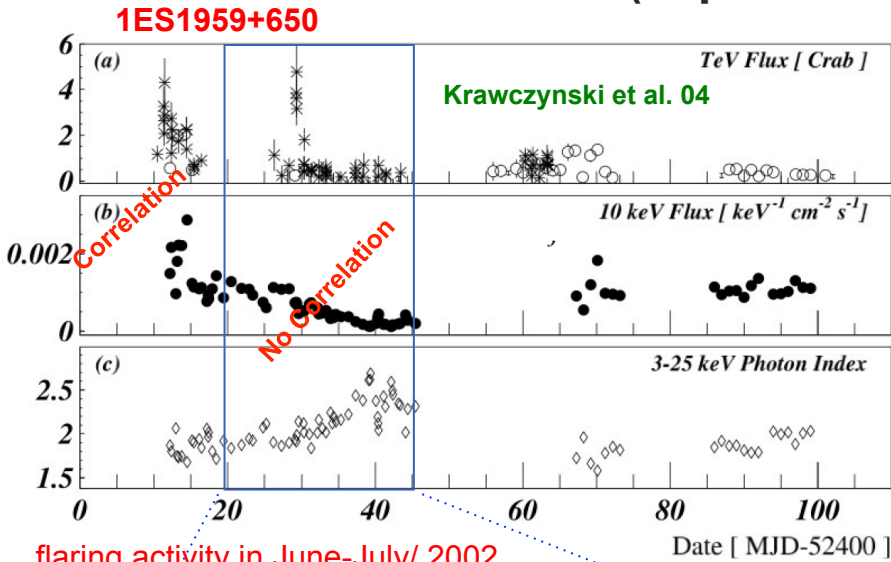
X-ray and TeV are correlated during the flare (flaring activity April/13).

Mrk421

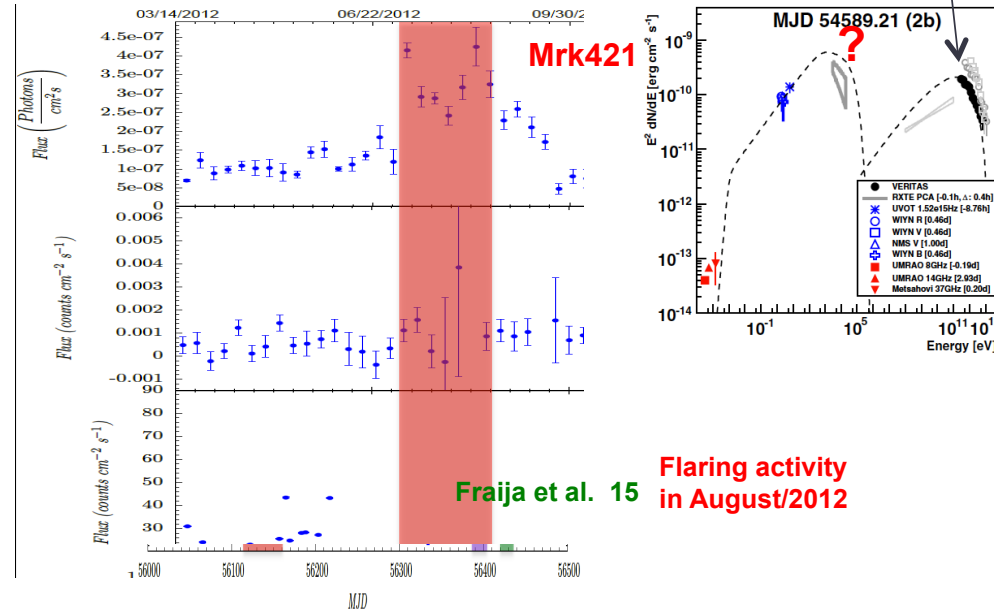
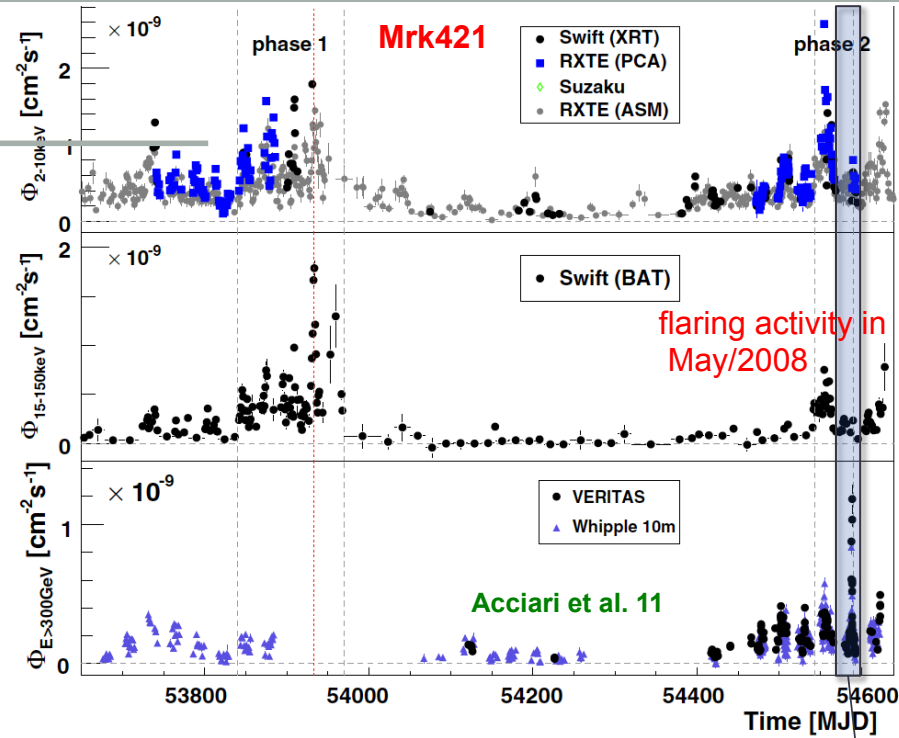
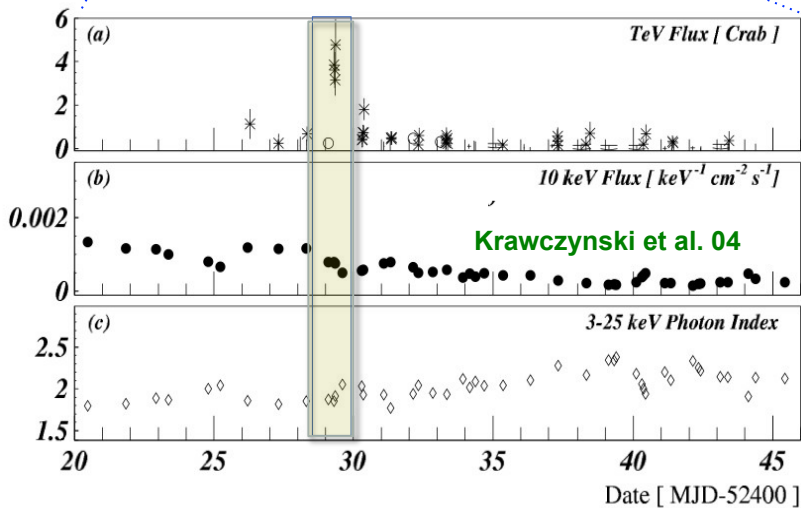
Mrk421

➤ Disconnections

(Orphan Flares)



flaring activity in June-July/ 2002



Is there a correlation between X-ray
and TeV γ -ray emission ????
(hourly, daily, monthly, etc.)

If so, what is the origin ??

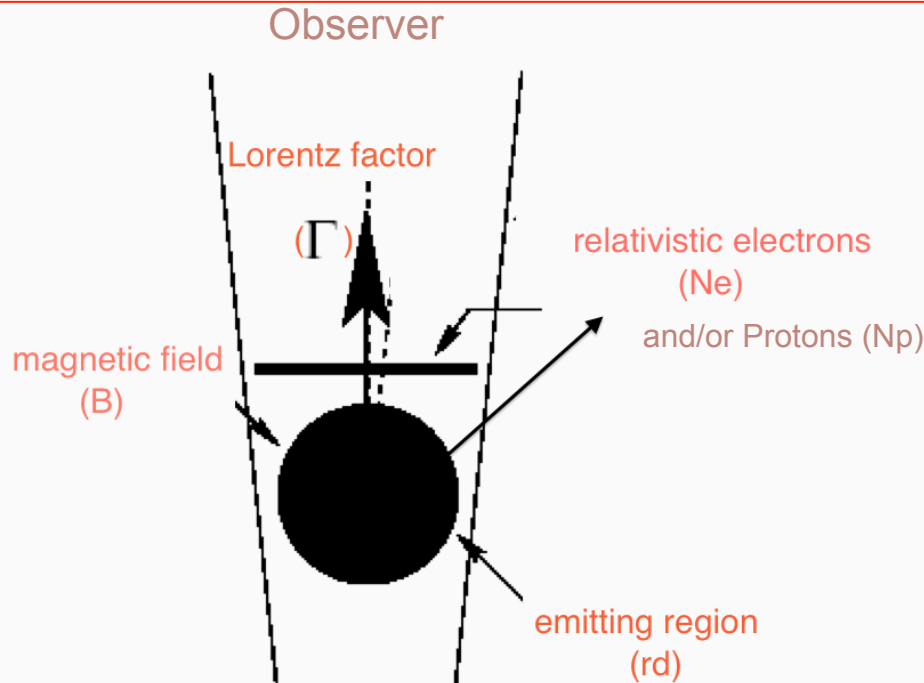
How could the orphan flares be explained ???

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➤ Theoretical Model

Sketch of the basic model



Material from the accretion disk is accreted to the BH and after is launched in the jet.

BH

Dermer et al. 03

We consider a spherical emitting region:

- Moving at relativistic speed with bulk Lorentz factor Γ .
- with a uniform particle densities (N_e and/or N_p).
- with radius (rd).
- Endowed with a magnetic field B .

Leptonic model

Just 4 parameters
(B , Γ , rd and $N_e(\beta)$)

Hadronic model

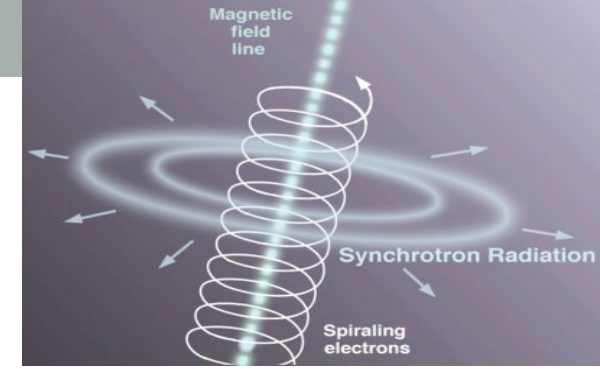
Additionally N_p
(5 parameters)

(we will use natural unities $c=1$ and prime quantities are in the comoving frame)

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➤ Synchrotron radiation



By considering that a fraction of total energy is given

To accelerate electrons $U_e = m_e \int \gamma_e N_e(\gamma_e) d\gamma_e$ \longrightarrow minimum e⁻ Lorentz factor $\gamma_{e,m} = \frac{(\alpha - 2)}{m_e(\alpha - 1)} \frac{U_e}{N_e}$

To generate and/or amplify magnetic field $U_B = \frac{B^2}{8\pi}$ \longrightarrow Cooling Time scale $t'_c = \frac{3m_e}{4\sigma_T} U_B^{-1} \gamma_e^{-1}$
 \downarrow cut-off e⁻ Lorentz factor $\gamma_{e,b} = \frac{3m_e}{4\sigma_T} (1 + Y)^{-1} \Gamma U_B^{-1} r_d^{-1}$

Compton parameter $Y = U_e / U_B$

Acceleration time scale $t'_{acc} = \sqrt{\frac{\pi}{2}} \frac{m_e}{q_e} U_B^{-1} \gamma_e$ \longrightarrow Maximum e⁻ Lorentz factor $\gamma_{e,max} = \left(\frac{9q_e^2}{8\pi\sigma_T} \right)^{1/4} U_B^{-1/4}$

Photon energy released $\epsilon_\gamma^{obs}(\gamma_e) = \sqrt{\frac{8\pi q_e^2}{m_e^2}} \Gamma U_B^{1/2} \gamma_{e,i}^2$

Break synchrotron energies

$$\begin{aligned} \epsilon_{\gamma,m}^{obs} &= \frac{\sqrt{8\pi} q_e (\alpha - 2)^2}{m_e^3 (\alpha - 1)^2} \Gamma U_B^{1/2} U_e^2 N_e^{-2}, \\ \epsilon_{\gamma,c}^{obs} &= \frac{9\sqrt{2\pi} q_e m_e}{8\sigma_T^2} (1 + Y)^{-2} \Gamma^3 U_B^{-3/2} r_d^{-2}, \\ \epsilon_{\gamma,max}^{obs} &= \frac{3q_e^2}{m_e \sigma_T} \Gamma, \end{aligned}$$

electron distribution

$$N_e(\gamma_e) \propto \begin{cases} \gamma_e^{-\alpha} & \gamma_{e,m} < \gamma_e < \gamma_{e,b}, \\ \gamma_{e,b} \gamma_e^{-(\alpha+1)} & \gamma_{e,b} \leq \gamma_e < \gamma_{e,max}, \end{cases}$$

Flux conservation

$$\epsilon_\gamma N_\gamma(\epsilon_\gamma) d\epsilon_\gamma \propto N_e(E_e) dE_e$$

Electron Synchrotron spectrum

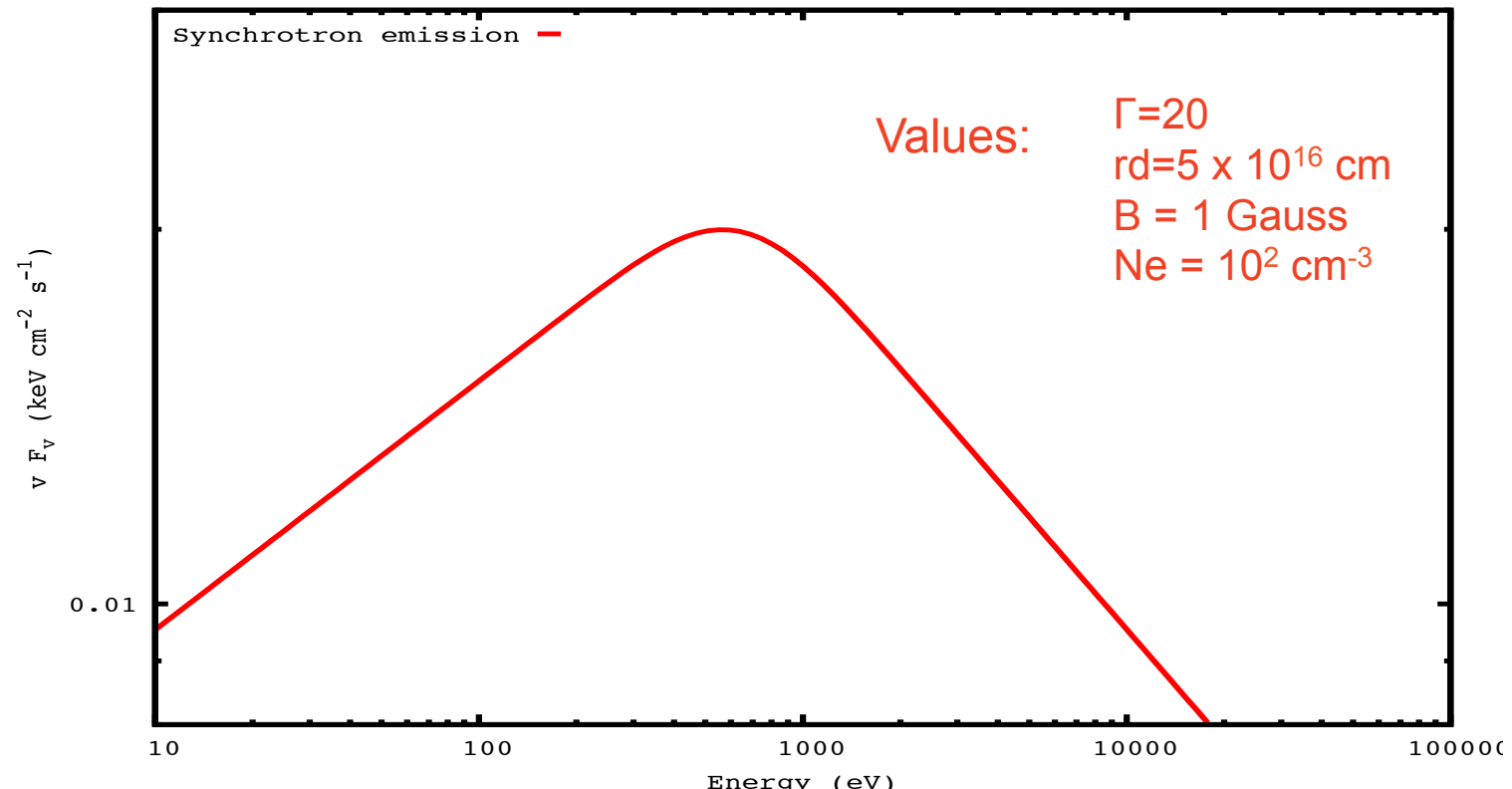
$$\epsilon_\gamma^2 N_\gamma(\epsilon_\gamma) = A_{syn,\gamma} \begin{cases} (\frac{\epsilon_\gamma}{\epsilon_{\gamma,m}})^{4/3} & \epsilon_\gamma^{obs} < \epsilon_{\gamma,m}^{obs}, \\ (\frac{\epsilon_\gamma}{\epsilon_{\gamma,m}})^{-(\alpha-3)/2} & \epsilon_{\gamma,m}^{obs} < \epsilon_\gamma^{obs} < \epsilon_{\gamma,c}^{obs}, \\ (\frac{\epsilon_{\gamma,c}}{\epsilon_{\gamma,m}})^{-(\alpha-3)/2} (\frac{\epsilon_\gamma}{\epsilon_{\gamma,c}})^{-(\alpha-2)/2}, & \epsilon_{\gamma,c}^{obs} < \epsilon_\gamma^{obs} < \epsilon_{\gamma,max}^{obs} \end{cases}$$



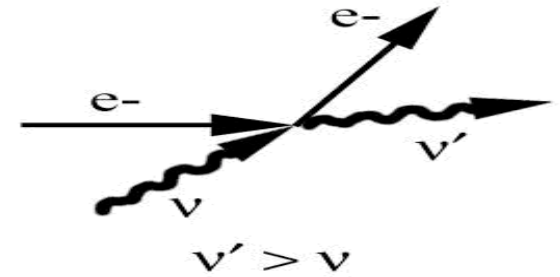
$$\beta = (\alpha + 2) / 2$$

$$A_{syn,\gamma} = \frac{P_{\nu,max}^{obs} n_e}{4\pi d_z^2} \epsilon_{\gamma,c}^{obs} \simeq \frac{m_e^2}{4q_e \sigma_T} d_z^{-2} \Gamma^5 U_B^{-1} N_e r_d$$

Milagro and Veritas (Low state)
Power index: $\beta = 2.3$



Compton scattering



High energy e- initially
e- loses energy

Synchrotron self-Compton relation

$$\epsilon_{\gamma, (m|c)}^{ic} \simeq \gamma_{e, (m|c)}^2 \epsilon_{\gamma, (m|c)}$$

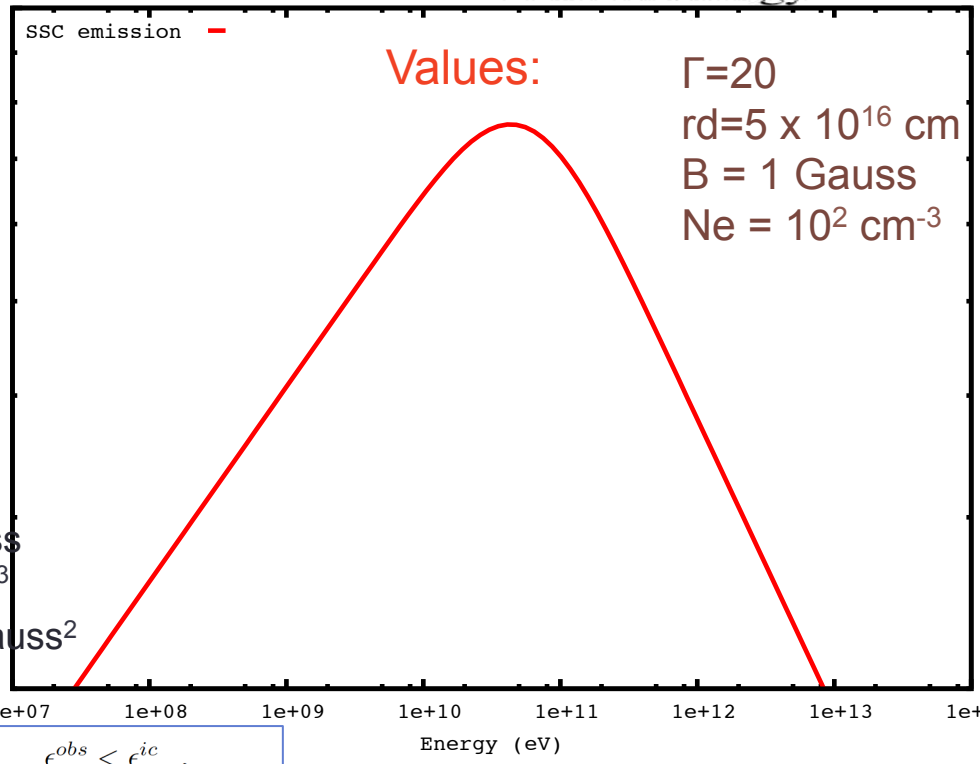
Break Compton scattering energies

$$\epsilon_{\gamma, m}^{ic} = \frac{\sqrt{8\pi} q_e (\alpha - 2)^4}{m_e^5 (\alpha - 1)^4} \Gamma U_B^{1/2} U_e^4 N_e^{-4},$$

$$\epsilon_{\gamma, c}^{ic} = \frac{81 \sqrt{2\pi} q_e m_e^3}{128 \sigma_T^4} (1 + Y)^{-4} \Gamma^5 U_B^{-7/2} r_d^{-4},$$

$$\epsilon_{\gamma, max}^{ic} = \frac{9 q_e^3}{2 \sqrt{2\pi} m_e \sigma_T^2} \Gamma U_B^{-1/2},$$

v Fv (erg cm⁻² s⁻¹)



Compton parameter

$$Y = 2.06 \times 10^{-5} N_0 / B_0^2$$

$B = B_0$ Gauss
 $N_e = N_0$ cm³
erg/cm³ = Gauss²

Compton scattering spectrum

$$\epsilon_{\gamma}^{ic2} N_{\gamma}(\epsilon_{\gamma}^{ic}) = Y (\epsilon_{\gamma}^{ic2} N_{\gamma})_{max}^{syn} \begin{cases} (\frac{\epsilon_{\gamma}^{ic}}{\epsilon_{\gamma, m}^{ic}})^{4/3} & \epsilon_{\gamma}^{obs} < \epsilon_{\gamma, m}^{ic} \\ (\frac{\epsilon_{\gamma}^{ic}}{\epsilon_{\gamma, m}^{ic}})^{-(\alpha-3)/2} & \epsilon_{\gamma, m}^{obs} < \epsilon_{\gamma}^{ic} < \epsilon_{\gamma, c}^{ic} \\ (\frac{\epsilon_{\gamma, c}^{ic}}{\epsilon_{\gamma, m}^{ic}})^{-(\alpha-3)/2} (\frac{\epsilon_{\gamma}^{ic}}{\epsilon_{\gamma, c}^{ic}})^{-(\alpha-2)/2} & \epsilon_{\gamma, c}^{ic} < \epsilon_{\gamma}^{ic} < \epsilon_{\gamma, max}^{ic} \end{cases}$$

$$(\epsilon_{\gamma}^{ic2} N_{\gamma})_{max}^{syn} = A_{syn, \gamma} \left(\frac{9 m_e^4 (\alpha - 1)^2}{16 \sigma_T^2 (\alpha - 2)^2} \Gamma^2 U_B^{-2} U_e^{-2} N_e^2 r_d^{-2} \right)^{-(\alpha-3)/2}$$

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➤ Py interactions

The target photon density

$$n_\gamma \simeq \frac{d_z^2}{r_d^2 \epsilon_{\gamma, \text{pk}}} (\nu F_\nu)$$

Charged and neutral pion production channels

$$p\gamma \longrightarrow \begin{cases} n \pi^+ & \text{fraction } 1/3, \\ p \pi^0 & \text{fraction } 2/3, \end{cases}$$

Photo pion cooling time,

$$t'_{\pi^0}{}^{-1} = \frac{1}{2 \gamma_p^2} \int d\epsilon \sigma_\pi(\epsilon) \xi_{\pi^0} \epsilon \int dx x^{-2} \frac{dn_\gamma}{d\epsilon_\gamma}(\epsilon_\gamma = x),$$

Photo pion efficiency

$$f_{\pi^0} = \frac{t'_d}{t'_{\pi^0}}$$

Proton distribution

$$\frac{dN_p}{dE_p} = A_p \left(\frac{E_p}{\text{GeV}} \right)^{-\alpha}$$

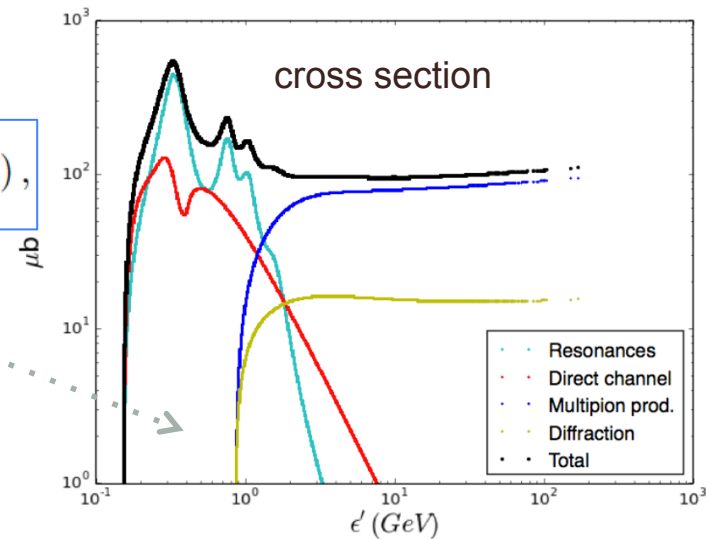
Photo pion spectrum:

$$\left(\epsilon^2 \frac{dN}{d\epsilon} \right)_{\pi^0, \gamma} = \frac{6 r_d \eta_\gamma \sigma_{\text{peak}} \Delta \epsilon_{\text{peak}}}{\Gamma^2 \epsilon_{\text{peak}} \left(\frac{2}{\xi_{\pi^0}} \right)^{\alpha-1} A_p} \begin{cases} \left(\frac{\epsilon_{\pi^0, \gamma, c}}{\epsilon_0} \right)^{-1} \left(\frac{\epsilon_{\pi^0, \gamma}}{\epsilon_0} \right)^{-\alpha+3} & \epsilon_{\pi^0, \gamma} < \epsilon_{\pi^0, \gamma, c} \\ \left(\frac{\epsilon_{\pi^0, \gamma}}{\epsilon_0} \right)^{-\alpha+2} & \epsilon_{\pi^0, \gamma, c} < \epsilon_{\pi^0, \gamma} \end{cases}$$

$$\pi^\pm \rightarrow \mu^\pm + \nu_\mu / \bar{\nu}_\mu$$

$$e^\pm + \nu_\mu / \bar{\nu}_\mu + \bar{\nu}_\mu / \nu_\mu + \nu_e / \bar{\nu}_e$$

$$\pi^0 \rightarrow \gamma\gamma$$



➤ μ^- , e^\pm and proton synchrotron radiation

Photon energy

$$\epsilon'_\gamma = \frac{3\pi q_e B'}{8 m_\mu^3} E_\mu'^2$$

where muon Lorentz factor

$$\gamma_i = m_i^2 / m_e^2 \gamma_e$$

Aharonian 2000, Mannheim 1993

Break synchrotron energies

$$\epsilon_{\gamma,c} = \frac{m_i^5}{m_e^5} \epsilon_{\gamma,c-e}$$

$$\epsilon_{\gamma,max} = \frac{m_i}{m_e} \epsilon_{\gamma,max-e}$$

Muons are accelerated by a power law distribution

$$N_i(\gamma_i): \gamma_i^{-\alpha} \text{ for } \gamma_i < \gamma_{i,b} \text{ and } \gamma_{i,b} \gamma_i^{-(\alpha+1)} \text{ for } \gamma_{i,b} \leq \gamma_i < \gamma_{i,max}$$

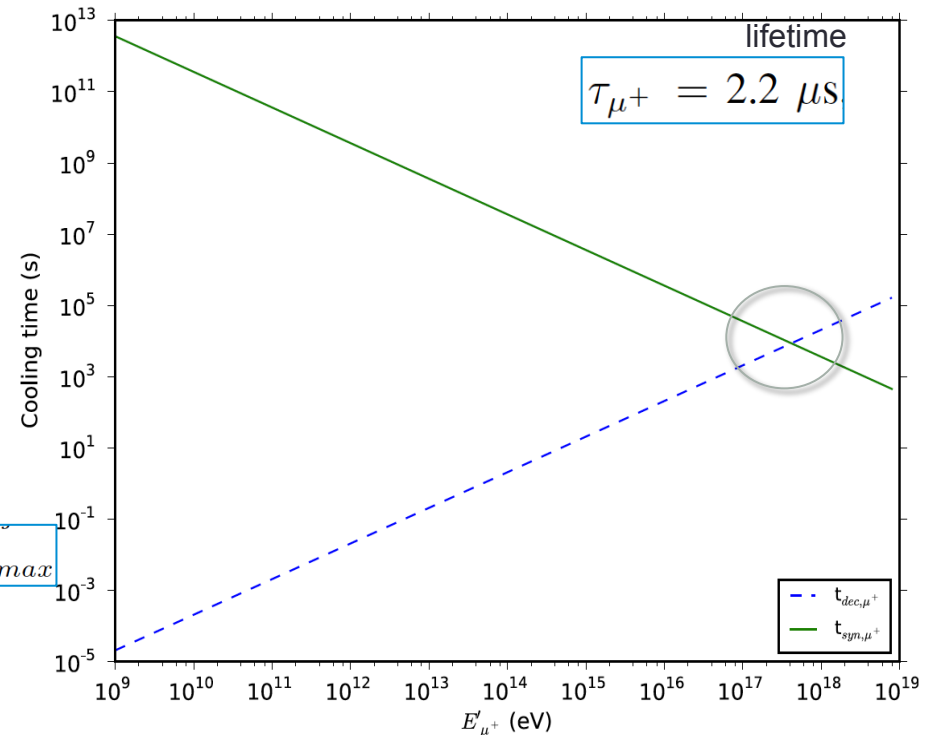
Synchrotron spectrum of muons and secondary pairs

$$\left(\epsilon^2 \frac{dN}{d\epsilon} \right)_{syn,\gamma} = \frac{4 \sigma_T m_i^4}{27 \pi^2 q_e m_e^3} r_j^3 D_z^{-2} \Gamma_j \epsilon_{\gamma,c-e} B' N_i \begin{cases} \left(\frac{\epsilon_{\gamma,c}}{\epsilon_0} \right)^{-1/2} \left(\frac{\epsilon_\gamma}{\epsilon_0} \right)^{-(\alpha-3)/2} & \epsilon_\gamma < \epsilon_{\gamma,c}, \\ \left(\frac{\epsilon_\gamma}{\epsilon_0} \right)^{-(\alpha-2)/2} & \epsilon_{\gamma,c} < \epsilon_\gamma < \epsilon_{\gamma,max} \end{cases}$$

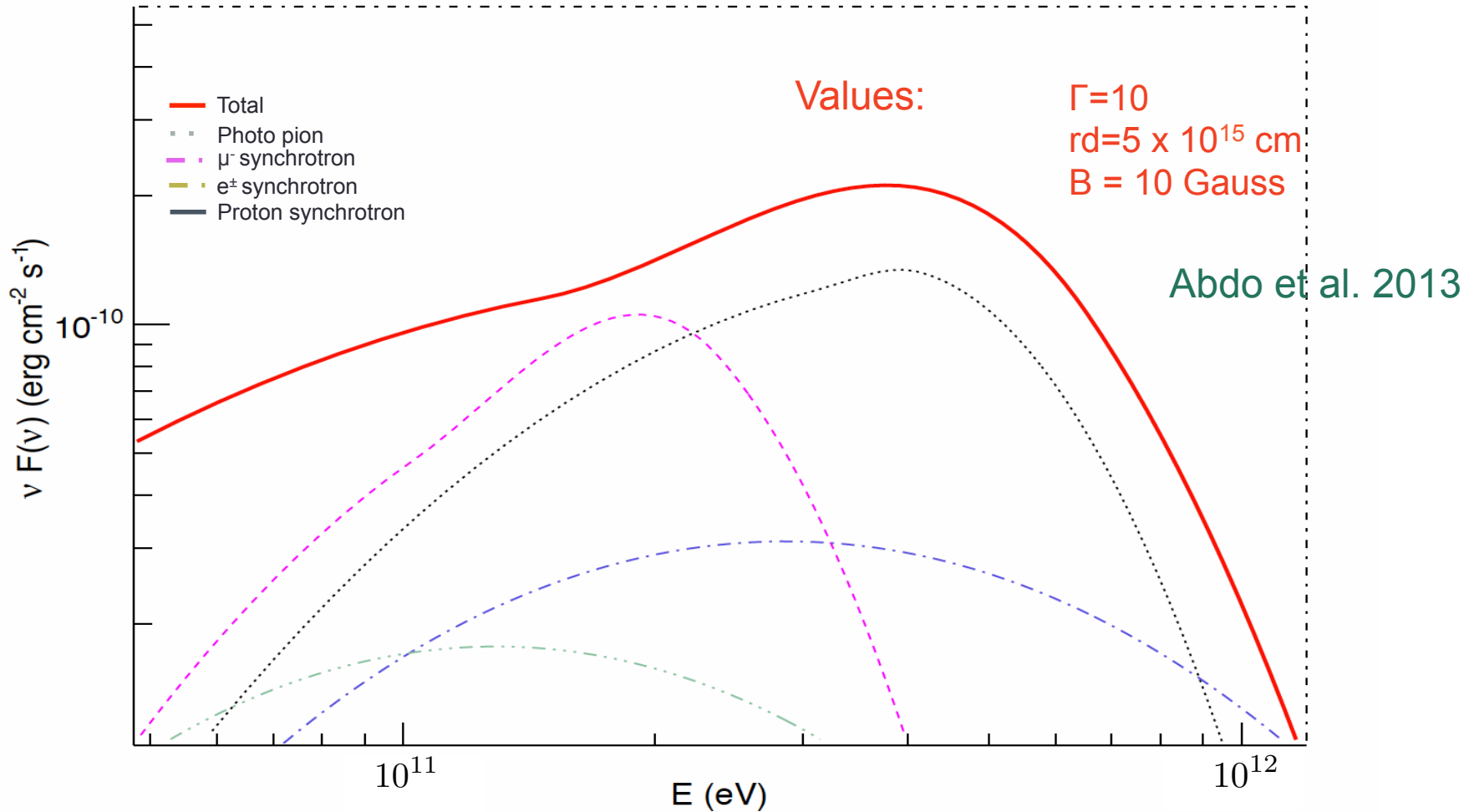
comparison

$$t'_{syn,i} = 6\pi m_i^4 / (\sigma_T m_e^2 B'^2 E_i')$$

$$t'_{\mu^+,dec} = \frac{E'_{\mu^+}}{m_{\mu^+}} \tau_{\mu^+}$$

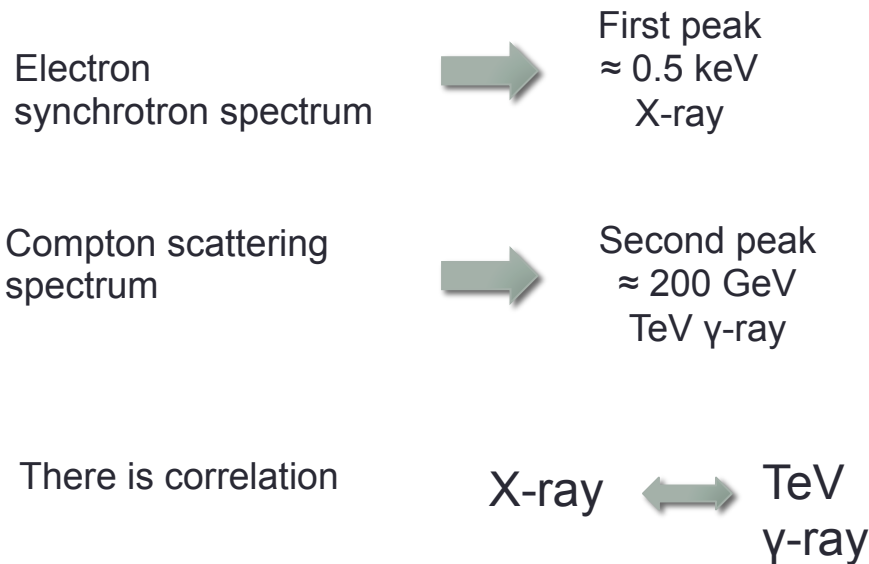


➤ Summing up: the hadronic spectrum

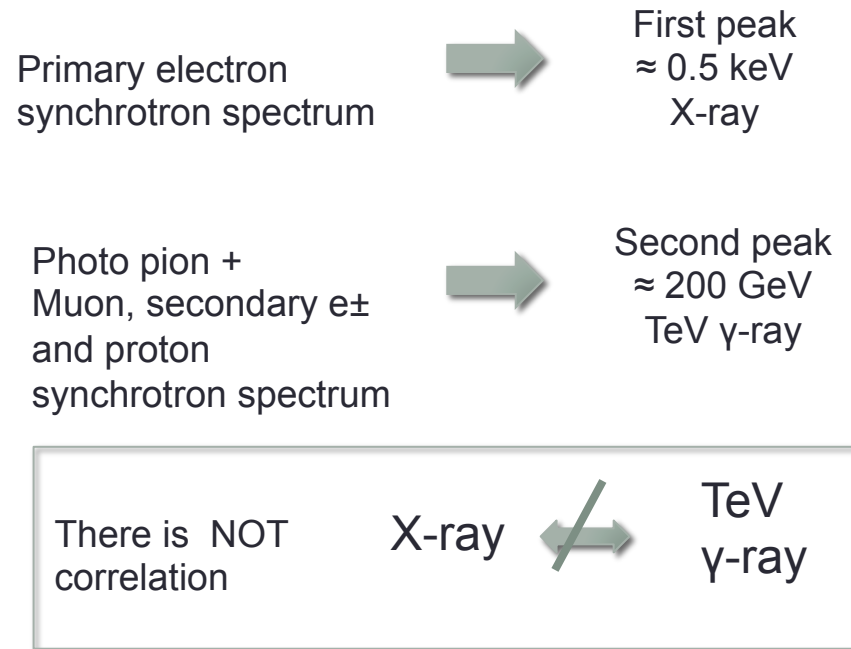


➤ Models

Leptonic



Hadronic



Orphan flares could be explain!!

Botcher 2004

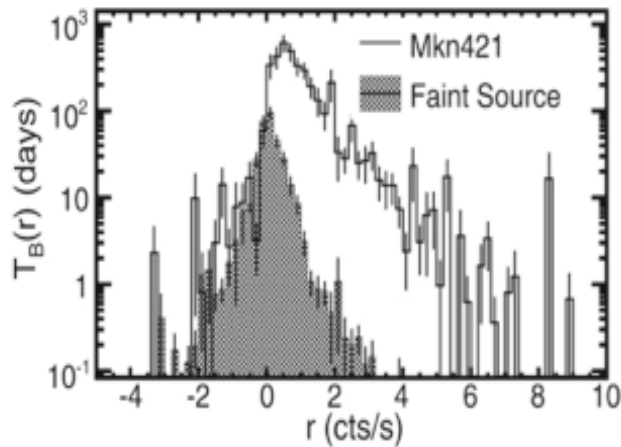
Fraija 2014

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➤ Data sets

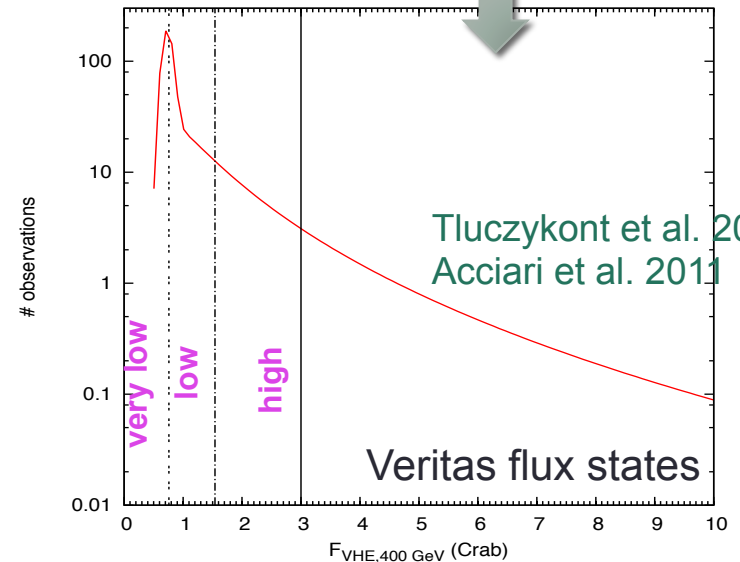
- Whipple - RXTE/ASM data (Acciari et al. 2014)
- Milagro - RXTE/ASM data
- HEGRA CT1 - RXTE/ASM (Aharonian et al. 2003)
- MAGIC/Whipple/VERITAS-XMM Newton data (Acciari et al. 2009)
- MAGIC-RXTE/ASM data (Albert et al. 2007)



Resconi et al. 2009

RXTE/ASM >14 years
2-10 keV
Mean=0.5 counts/s
(distribution)

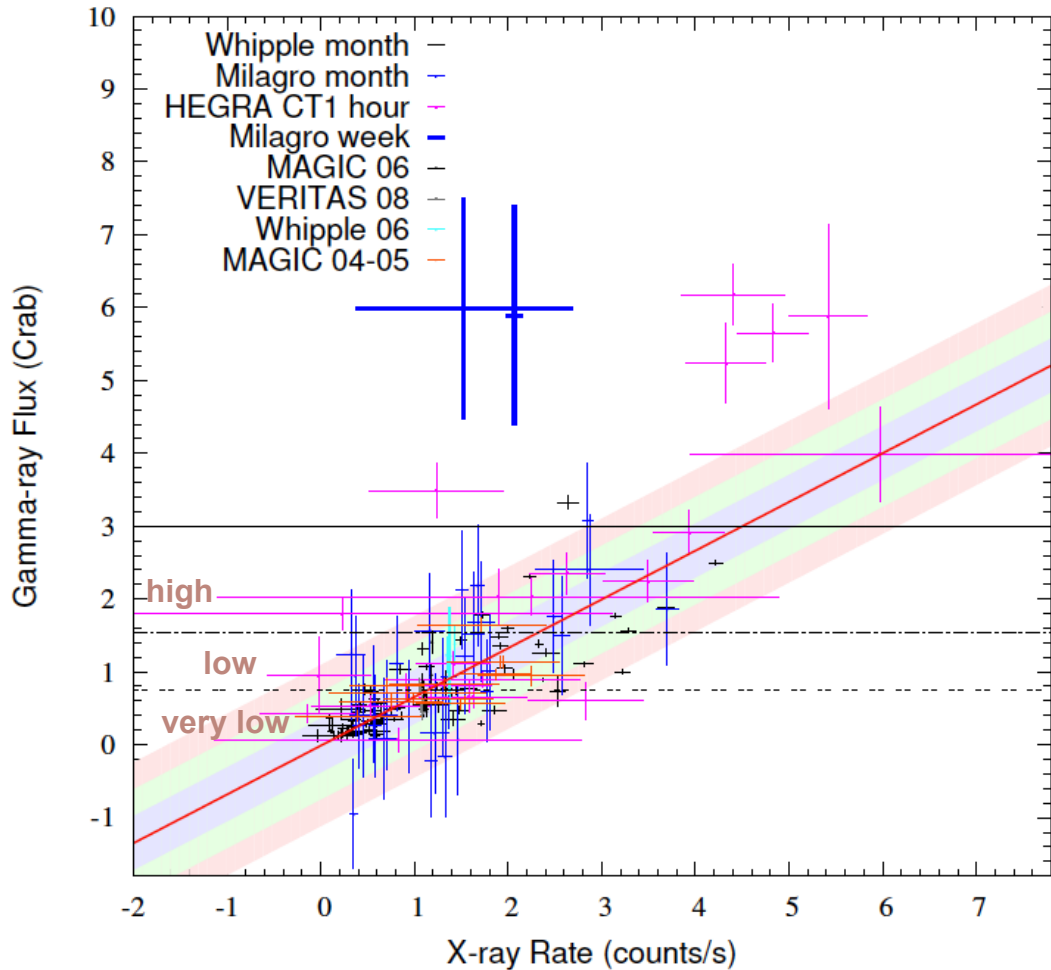
IACTs: MAGIC, H.E.S.S.,
Whipple/VERITAS, CAT, HEGRA
14 years $E > 400$ GeV
Mean=0.73 Crab
(distribution)



Gluczykont et al. 2010
Acciari et al. 2011

Veritas flux states

$F_{VHE,400 \text{ GeV}}$ (Crab)



- $R=0.73$
- Linear fit:
slope= 0.67 ± 0.04
- Sigma:
 $\sigma=0.37 \pm 0.03$

Patricelli et al 2015

➤ Considerations in our model

Values of parameters

$$\left\{ \begin{array}{l} \beta = 2.3 \\ r_d = 5 \times 10^{16} \text{ cm} \\ \Gamma = 10 \end{array} \right.$$

Acciari et al. 2011, 2014
Abdo et al. 2014

Values N_e and B are computed fitting data

VHE gamma-ray fluxes was normalized to the Crab flux, as measured by VERITAS



$$1 \text{ Crab} = 0.871 \times 10^{-10} \text{ erg/cm}^2/\text{s}$$

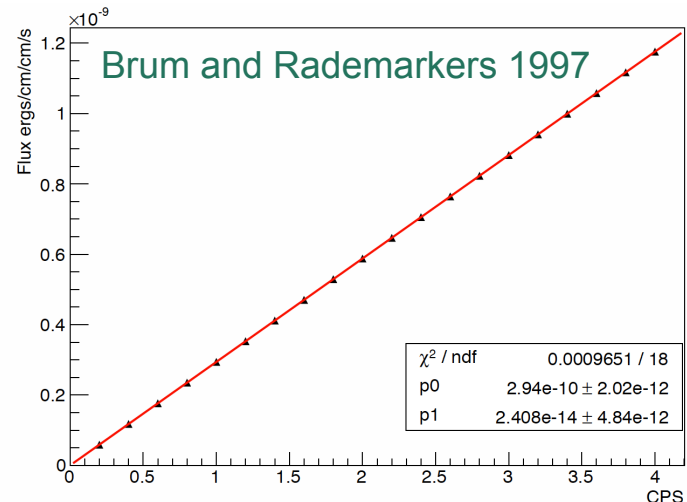
Unities of X-ray flux were changed (with the online WebPIMMS tool)

$$\text{erg/cm}^2/\text{s} \rightarrow \text{CPS}$$

http://heasarc.gsfc.nasa.gov/cgi-bin/Tools/w3pimms/w3pimms_pro.pl

Kalberla et al 2005

Column density fixed to the Galactic value in the direction of Mrk 421: $1.61 \times 10^{20} \text{ cm}^{-2}$



➤ Comparison

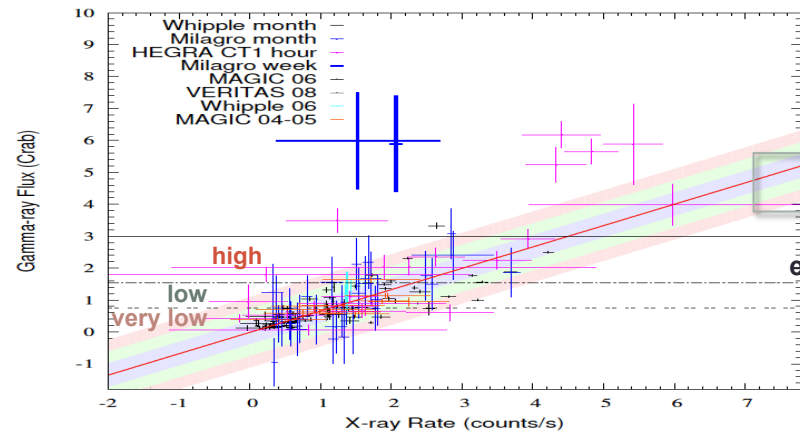
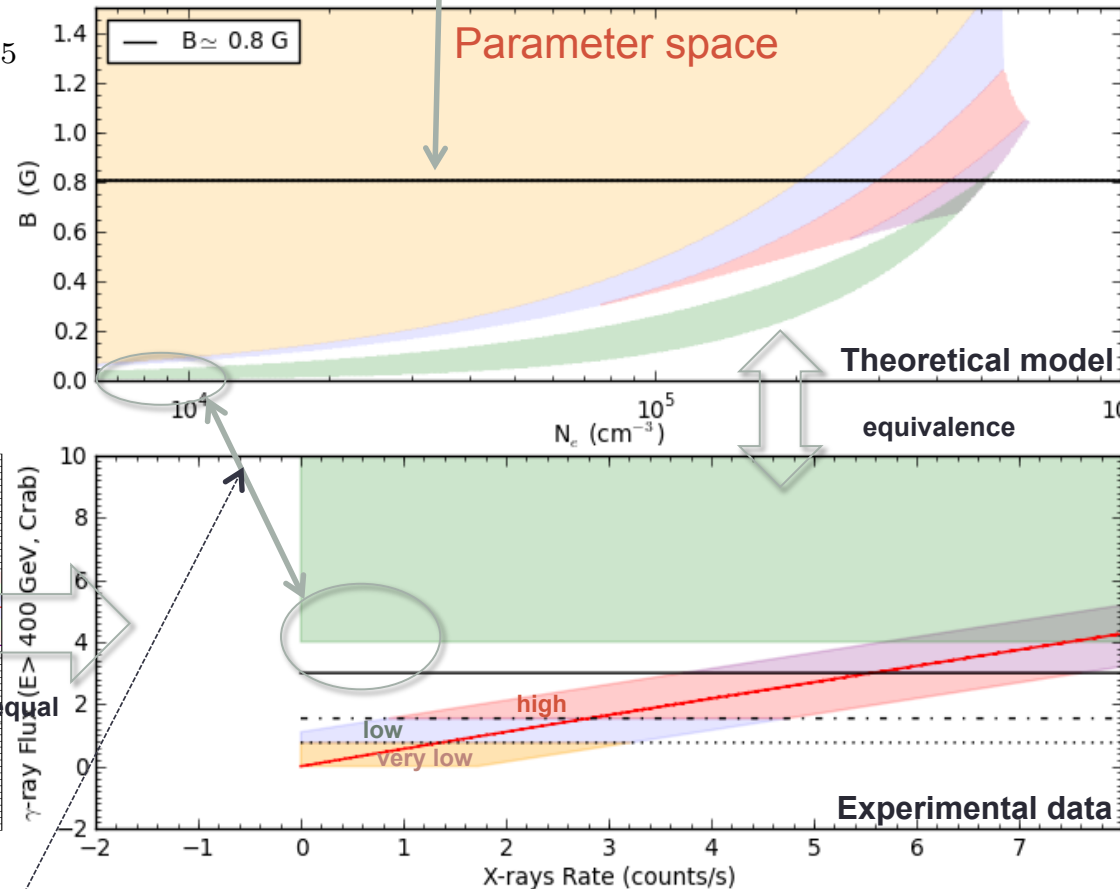
Magnetic field $0.01 \leq B \text{ (G)} \leq 1.6$

Electron density $10^3 \leq N_e \text{ (cm}^{-3}\text{)} \leq 10^{5.5}$

As fluxes increase the set of parameters decreases.

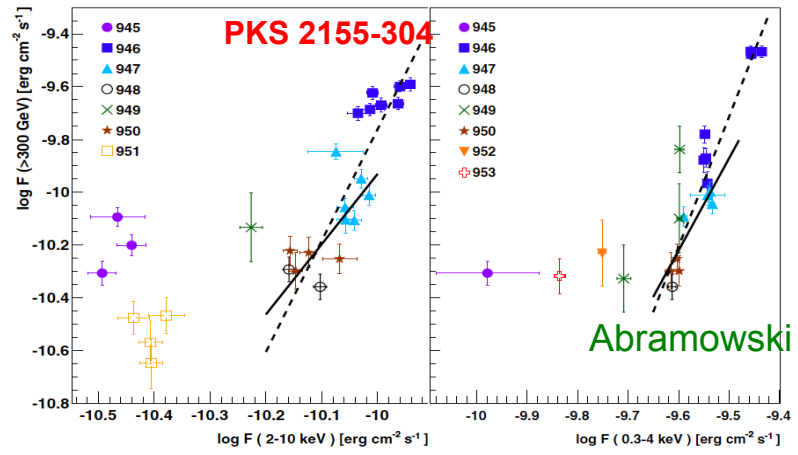
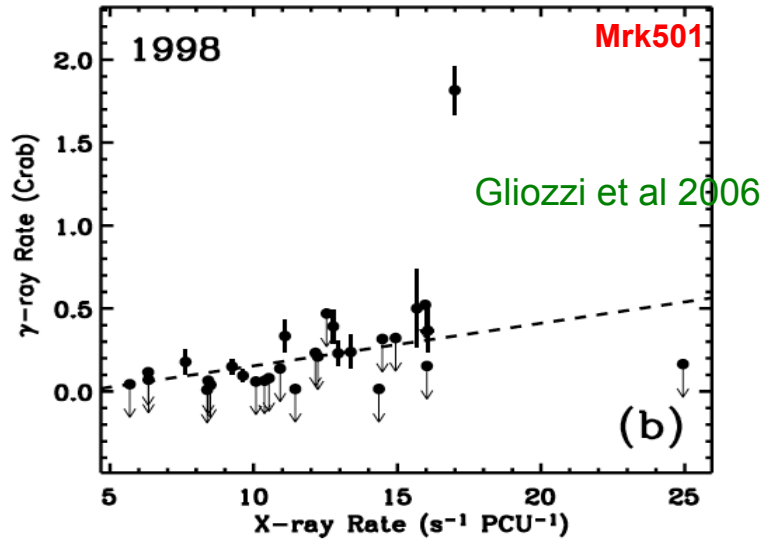
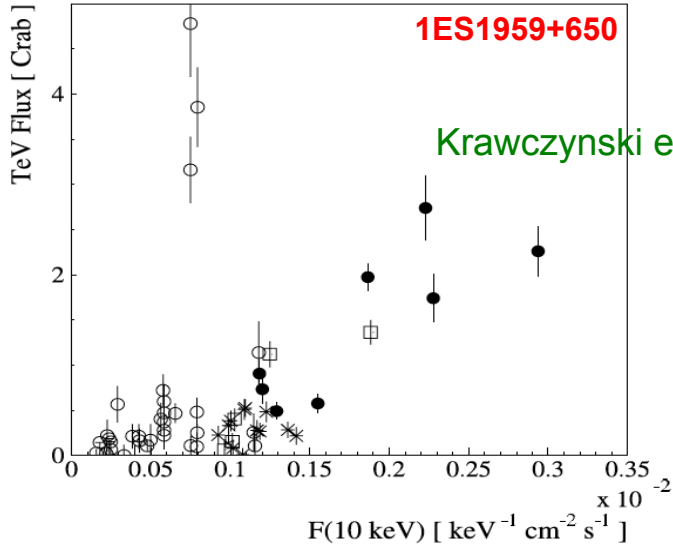
As fluxes decrease B decreases

Unique correlation $B \simeq 0.8 \text{ G}$



orphan flares

This model could be generalized to other blazars



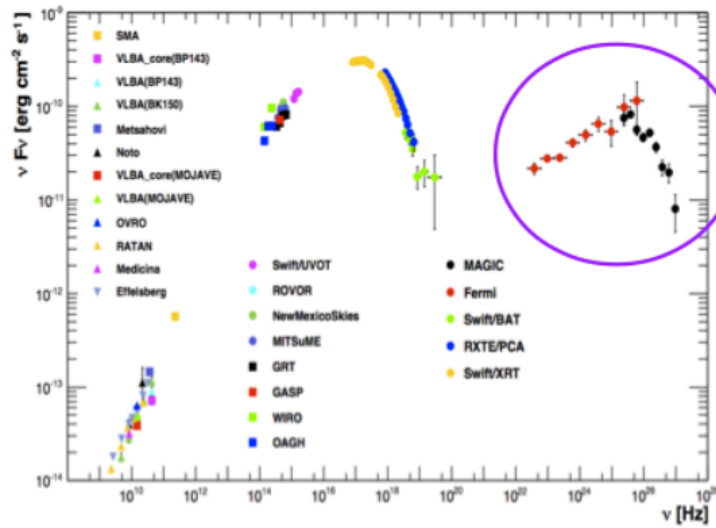
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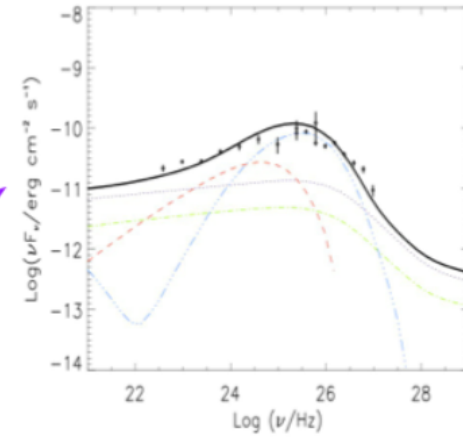
➤ Conclusions

- ✓ Mrk421 shows a correlation between TeV γ -ray and X-ray emissions independent of time scales and instruments, although it seems to break at the highest γ -ray fluxes.
- ✓ The overall correlation can be interpreted as SSC scenario with a single value of magnetic field $B \simeq 0.8$ G
- ✓ The “outliers” and orphan flares might be described within SSC framework for a set of values of N_e and B different from the ones characterizing the overall correlation.
- ✓ We have developed a theoretical model that can explain the correlation between TeV γ -ray and X-ray emissions of Mrk 421. Although the hadronic model can describe the SED, it can not explain the correlation.

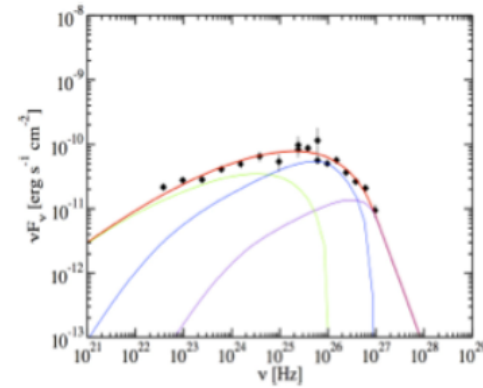
Back up



Hadronic



SSC – 3 electron populations



Both leptonic and hadronic models describe well the SED