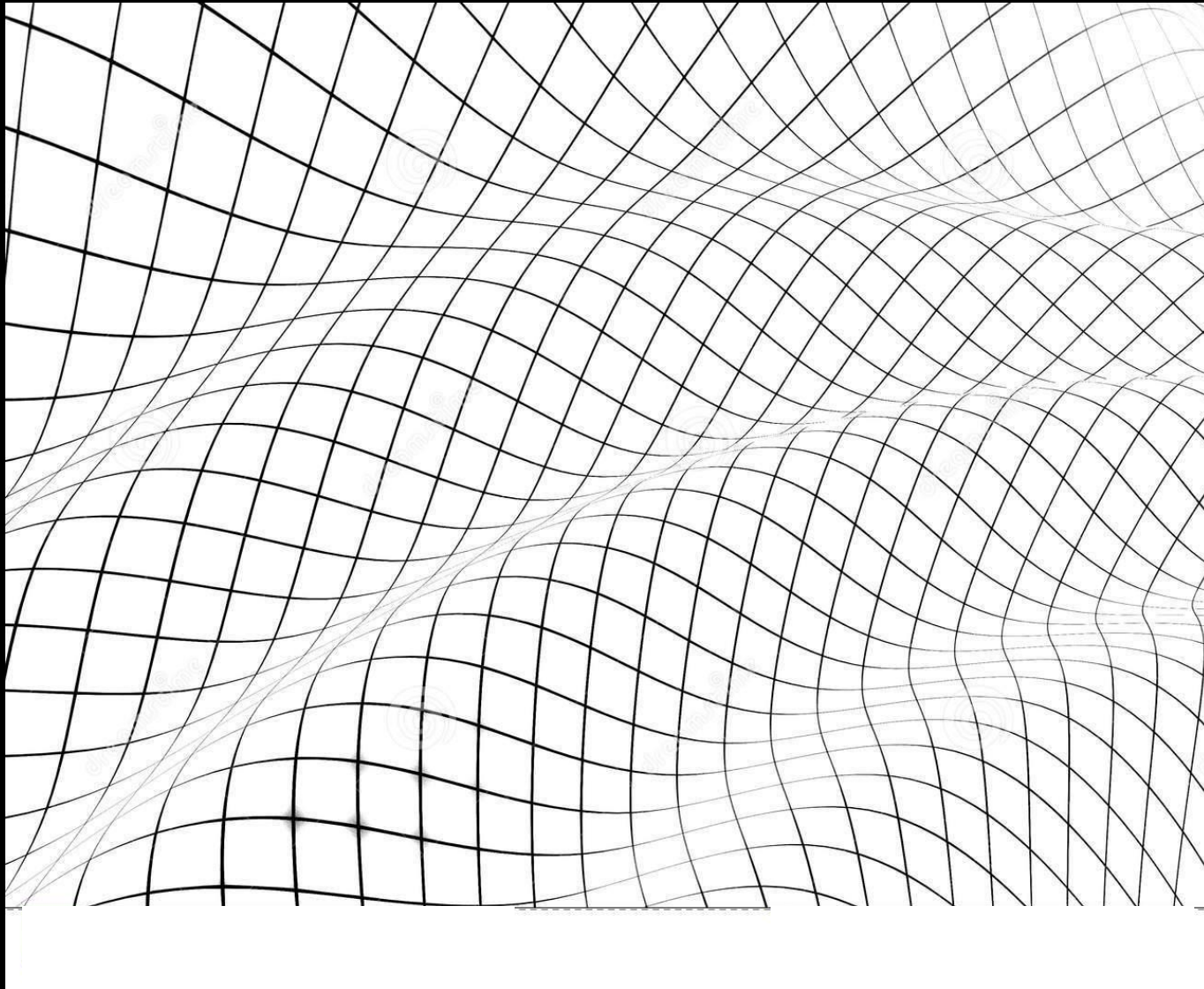


Light as a probe of the structure of space-time

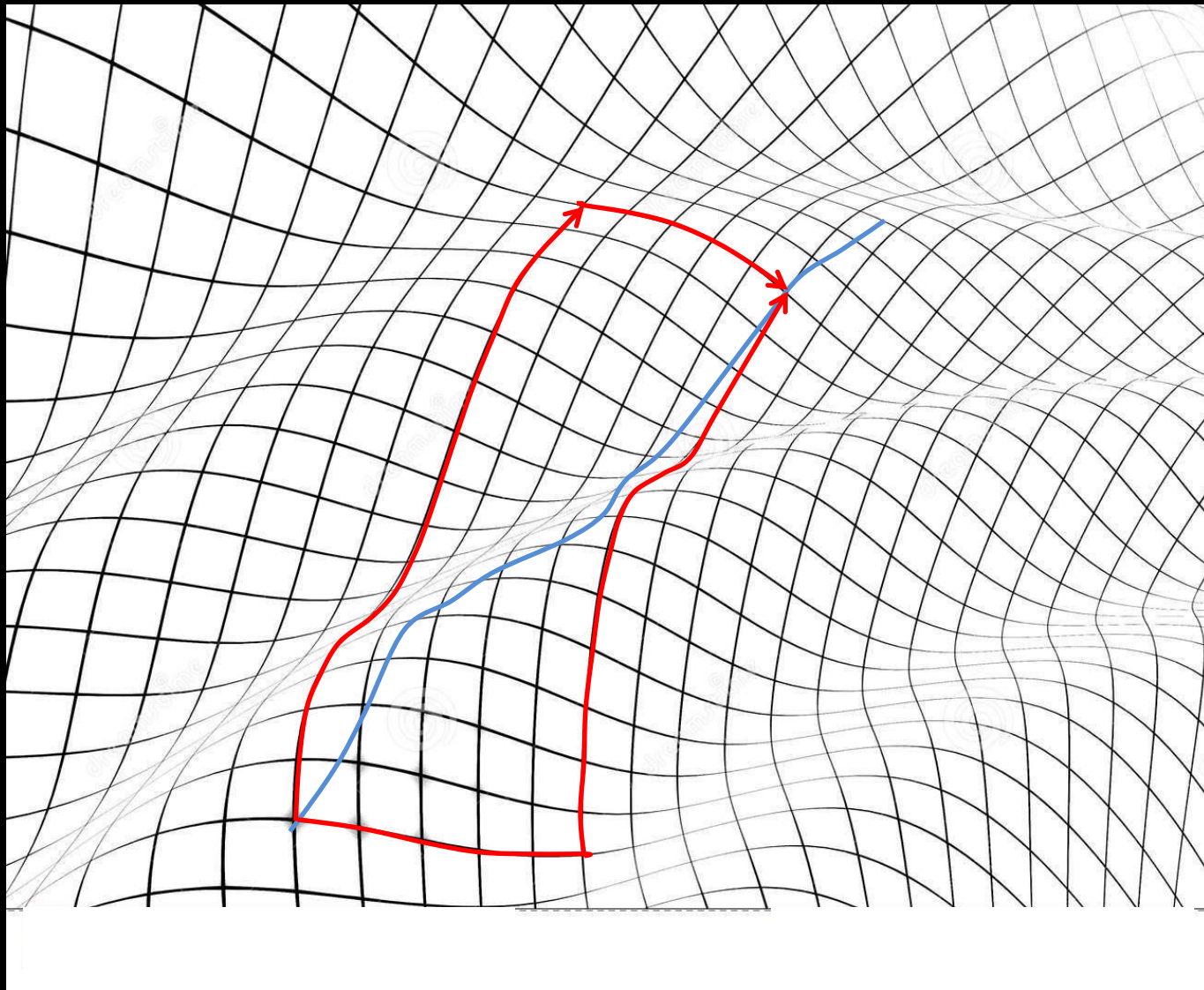
Angelo Tartaglia

Politecnico di Torino and INFN

Best covering of space-time by means of null geodesics



Best covering of space-time by means of null geodesics




Chiral symmetry with respect to time

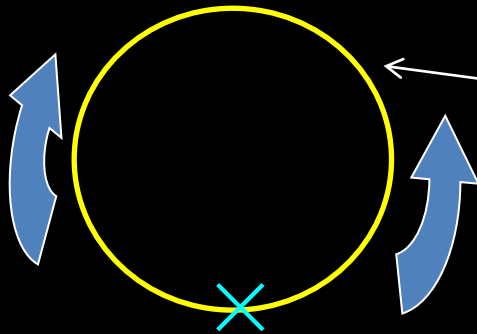
$$ds^2 = g_{00}c^2 dt^2 + 2g_{0i}cdtdx^i + g_{ij}dx^i dx^j$$

Light $\rightarrow ds = 0$

$$dt = \frac{-g_{0i}dx^i \left(\pm \sqrt{\left(g_{0i}dx^i\right)^2 - g_{00}g_{ij}dx^i dx^j} \right)}{cg_{00}}$$



Closed space trajectory



Counterrotating light beams

Shape is irrelevant

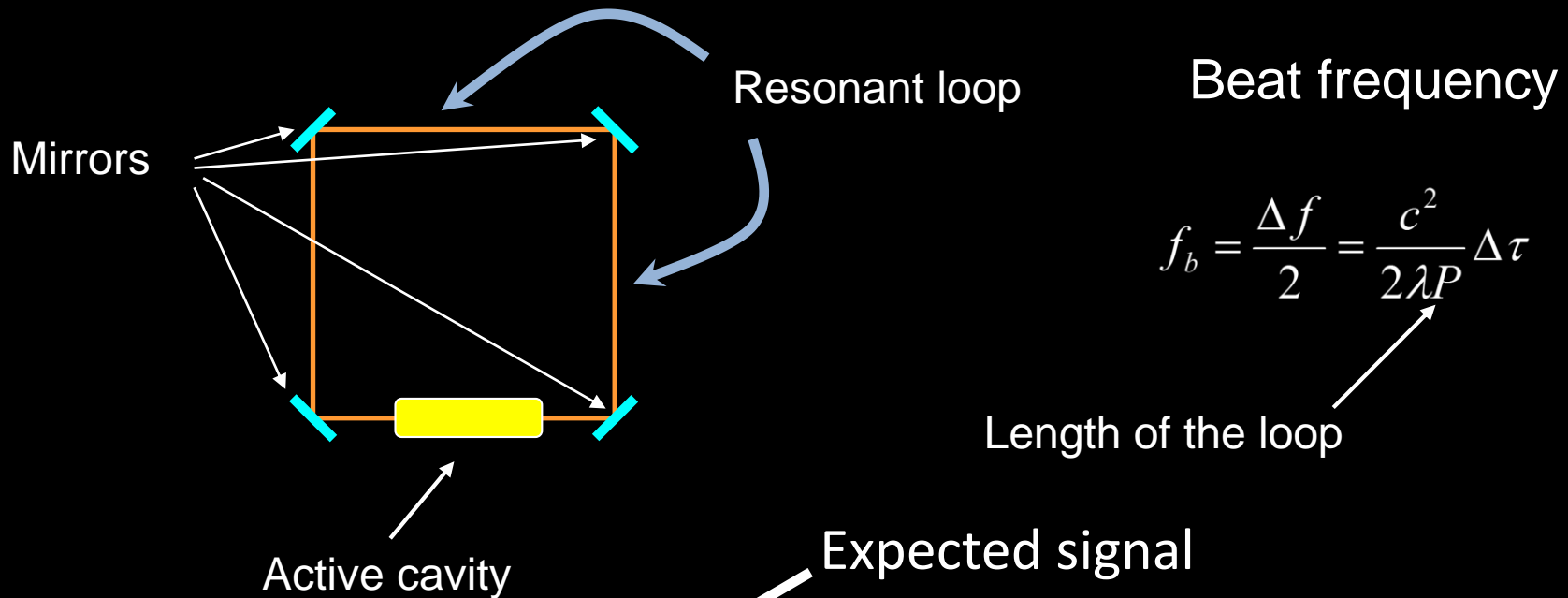
$$x^i = x^i(\ell)$$

$$dt = \frac{-g_{0i} \frac{dx^i}{d\ell} d\ell + \sqrt{\left(g_{0i} \frac{dx^i}{d\ell}\right)^2 - g_{00}g_{ij} \frac{dx^i}{d\ell} \frac{dx^j}{d\ell}} |d\ell|}{cg_{00}}$$

Time of flight
difference

$$\Delta\tau = \tau_+ - \tau_- = -\frac{2}{c} \sqrt{g_{00}} \oint \frac{g_{0i}}{g_{00}} \frac{dx^i}{d\ell} d\ell$$

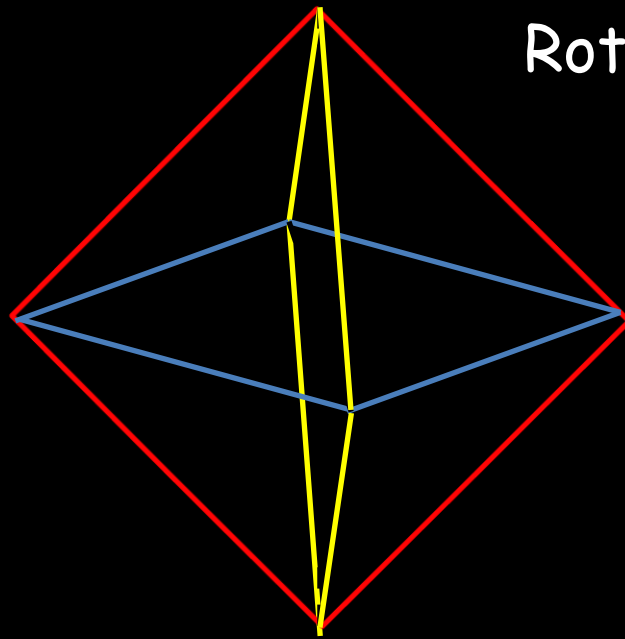
A ringlaser on earth



$$f_b = 2 \frac{A}{\lambda P} \left[\vec{\Omega} - 2G \frac{M}{c^2 R} \Omega \sin \theta \hat{u}_\theta + \frac{GJ_\oplus}{c^2 R^3} (2 \cos \theta \hat{u}_r + \sin \theta \hat{u}_\theta) \right] \cdot \hat{u}_n$$

GINGER

Underground at Gran Sasso



Rotation speed of the Earth

$$\Omega = 7.2 \times 10^{-5} \text{ s}^{-1}$$

GR terms

$$\Omega_G \approx 10^{-10} \Omega$$
$$\Omega_{LT} \approx 10^{-9} \Omega$$

Details in a poster presented at this conference

PPN parameters

$$ds^2 = \left(1 - 2\frac{GM}{c^2 r}\right) d\tau^2 - \left(1 + 2\gamma\frac{GM}{c^2 r}\right) dr^2 - r^2 d\vartheta^2$$

$$- r^2 \sin^2 \vartheta d\phi^2 + (4 + 4\gamma + \alpha_1) \frac{GJ}{2c^3 r} \sin^2 \vartheta d\tau d\phi$$

$$2M \rightarrow (1 + \gamma)M \qquad 2J \rightarrow \left(1 + \gamma + \frac{\alpha_1}{4}\right)J$$

Changes in the GR effects induced by ppn terms

New expected results including the effects of γ and α_1

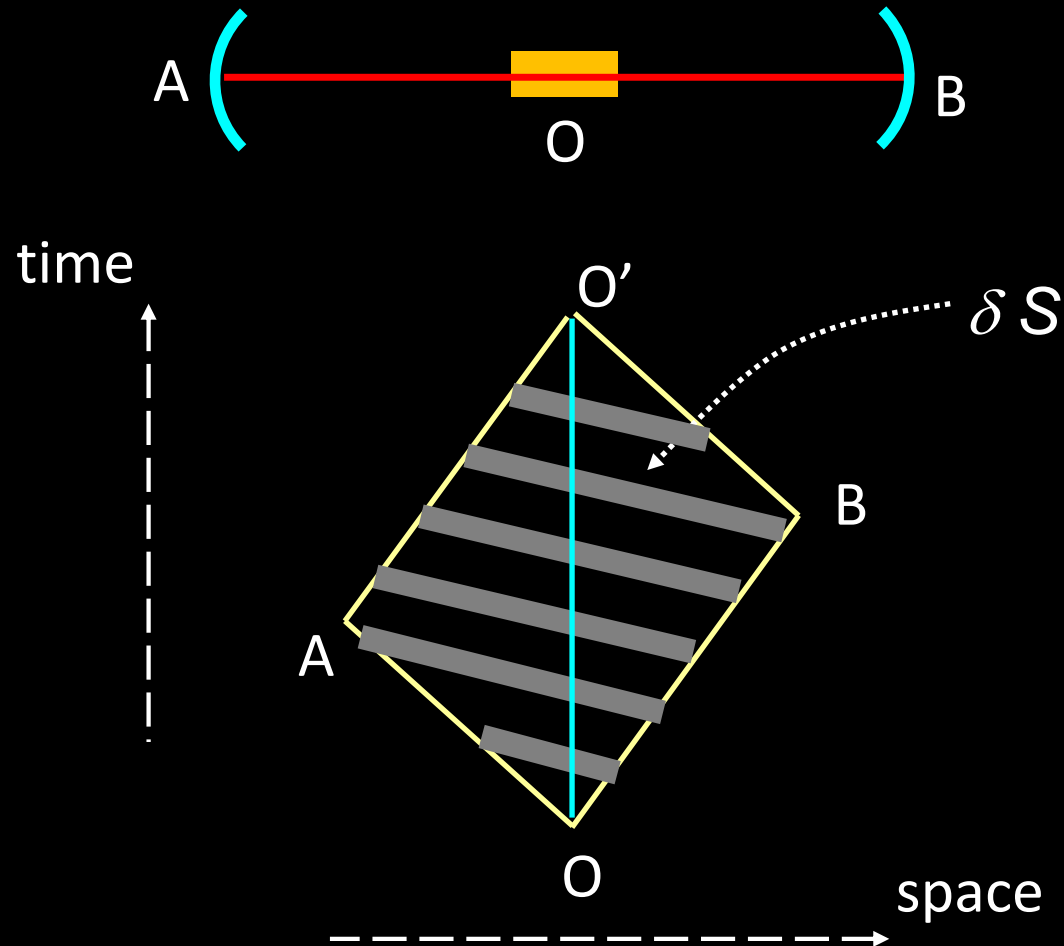
$$\Omega_G = -[2 + (\gamma - 1)]G \frac{M}{c^2 R} \Omega \sin \theta \hat{u}_\theta \cdot \hat{u}_n$$

$$\Omega_{LT} = \frac{2 + (\gamma - 1) + \alpha_1 / 4}{2} G \frac{J}{c^2 R^3} (2 \cos \theta \hat{u}_r + \sin \theta \hat{u}_\theta) \cdot \hat{u}_n$$

GR: $\gamma = 1$ $\alpha_1 = 0$

Present constraints on the parameters impose accuracies better than 10^{-22} rad/s.

A linear cavity (and GW)



The effect of curvature and anisotropy

$$\delta F^{\mu\nu} = \left(R^{\mu}_{\ \varepsilon 0 i} F^{\varepsilon\nu} + R^{\nu}_{\ \varepsilon 0 i} F^{\mu\varepsilon} \right) \delta S^{0i}$$

Electromagnetic tensor

Riemann tensor

Space-time area
spanned by the cavity

Depends on the orientation of the cavity

Approximated Riemann tensor.

An example:

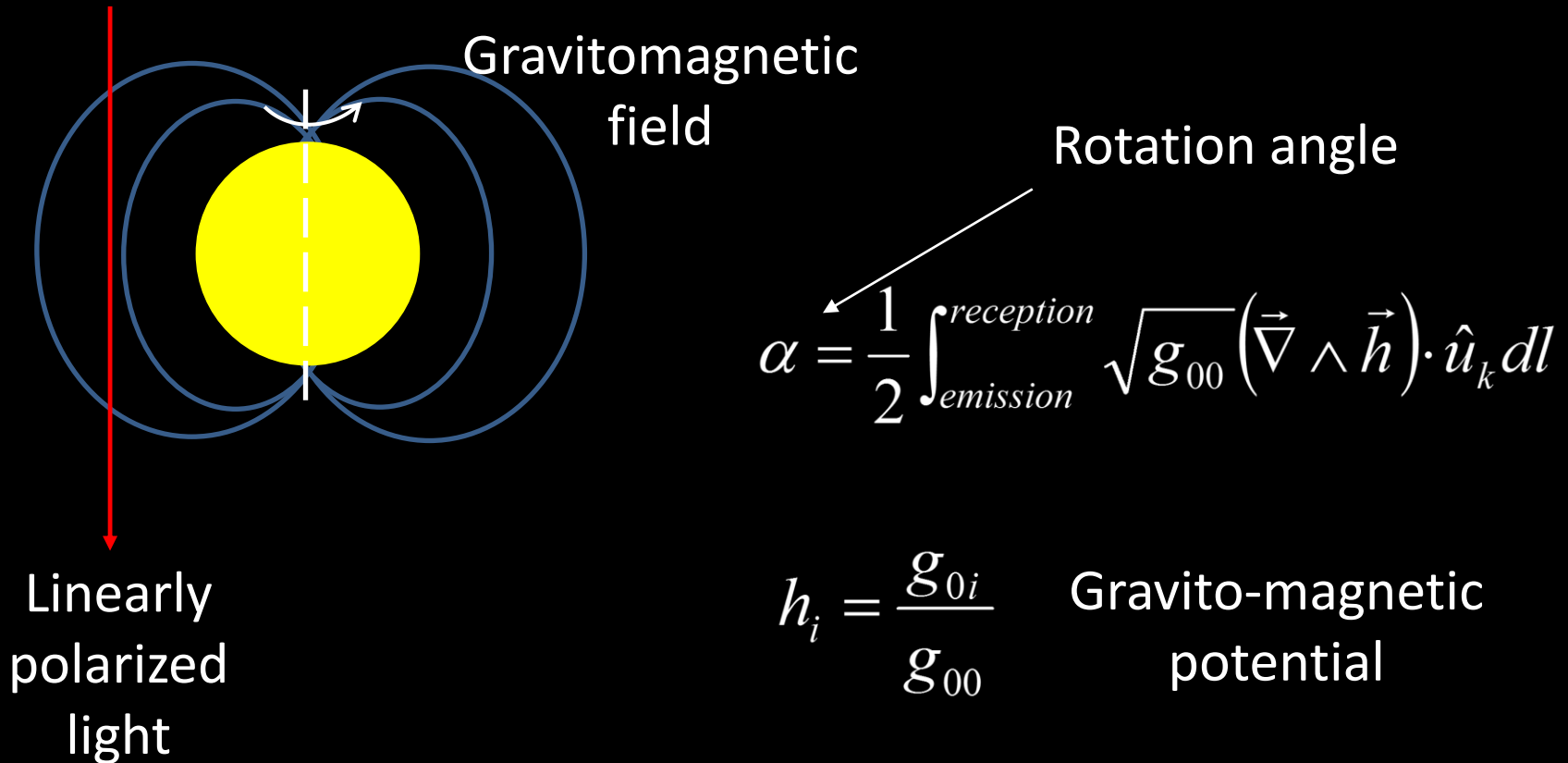
$$R^{\phi}_{r0\theta} \delta F^{\vartheta\phi} \cong \left[\left(\frac{(GM)^{3/2}}{c^3 R^{7/2}} - 3 \frac{GJ}{c^3 R^4} \right) \frac{\cos \vartheta}{\sin^2 \vartheta} \right] \frac{l^2}{R} F^{\vartheta r}$$

Length of the cavity

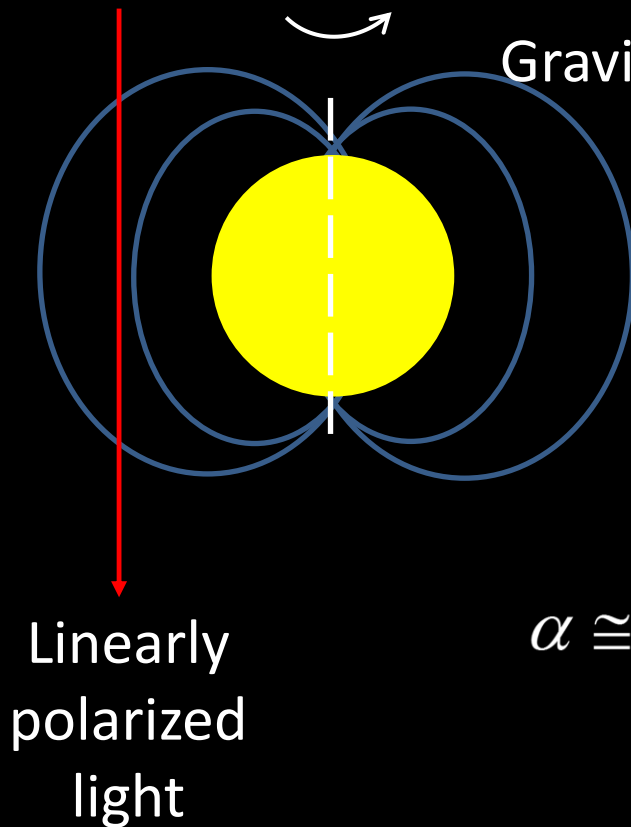
Change in the radial component of the magnetic field

East-West component of the magnetic field

Gravito-magnetic Faraday effect



Weak field approximation (lowest order in the angular momentum)



Approximately straight propagation parallel to the axis of the central mass

$$\alpha \cong \int_{-\infty}^{+\infty} (\vec{\nabla} \wedge \vec{g}_0) \cdot \hat{u}_k dl = \frac{j}{4b^2}$$

\swarrow GJ/c^2
 \nwarrow Minimum approach distance

Conclusion

- Light is an intrinsically relativistic probe of space-time
- Closed paths of light in space (and space-time) evidence general relativistic effects either by
 - Interferometric techniques
or
 - Beat frequency measurements (ring lasers \leftrightarrow GINGER)
- Improved observation techniques may give access to gravito-magnetic effects on the polarization of light



Collaboration

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