

A combined limit for neutrinoless double-beta decay

Pawel Guzowski

PHYSICAL REVIEW D **92**, 012002 (2015)

Combined limit on the neutrino mass from neutrinoless double- β decay and constraints on sterile Majorana neutrinos

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The University of Manchester

Outline

- Combined limit from 6 experimental inputs:
 - CUORICINO and CUORE-o
 - EXO-200
 - GERDA
 - KAMLAND-Zen
 - NEMO-3
- Cross-check of these individual results
- The combination
 - Nuclear Matrix Element model dependence
- Limits on sterile Majorana neutrinos

Introduction

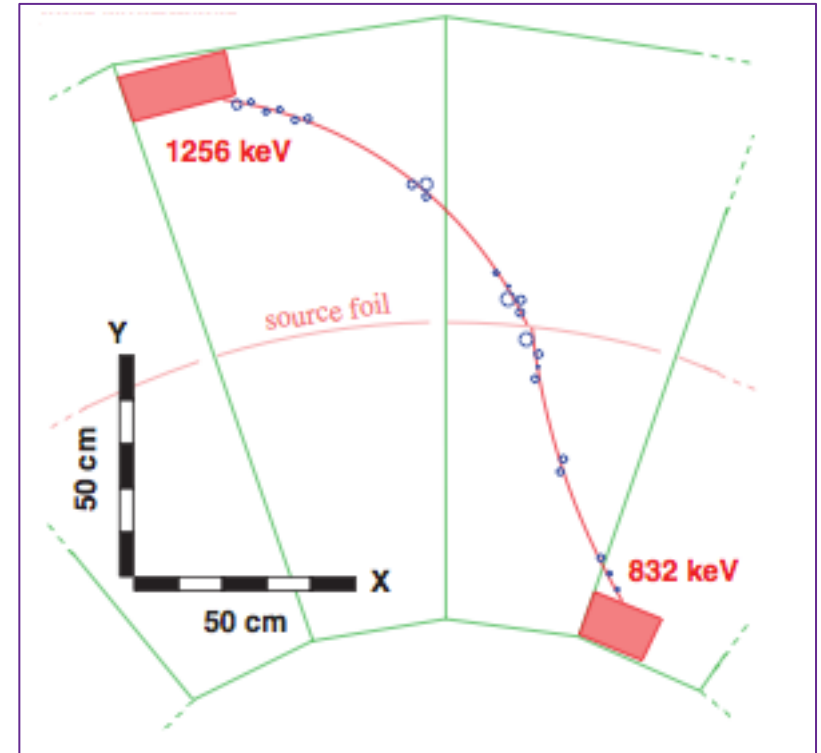
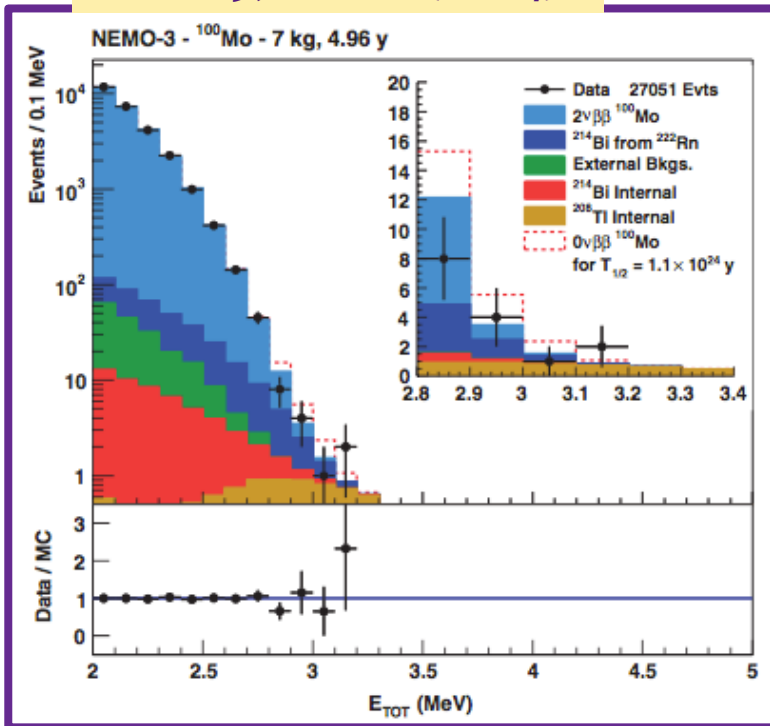
- Limits on $0\nu\beta\beta$ have only so far been set for single experiments, or for a single isotope
- Combining results for multiple isotope is important, for improving sensitivity, and also for disentangling mechanisms if an observation is made
- We introduce a method for combining results of different experiments, and apply it to 6 recent experimental results
 - COURICINO/CUORE-0, EXO-200, GERDA, KamLAND-Zen & NEMO-3

THE EXPERIMENTS

NEMO-3

- Tracker-calorimeter
 - ^{100}Mo isotope

PRD 89, 111101 (2014)



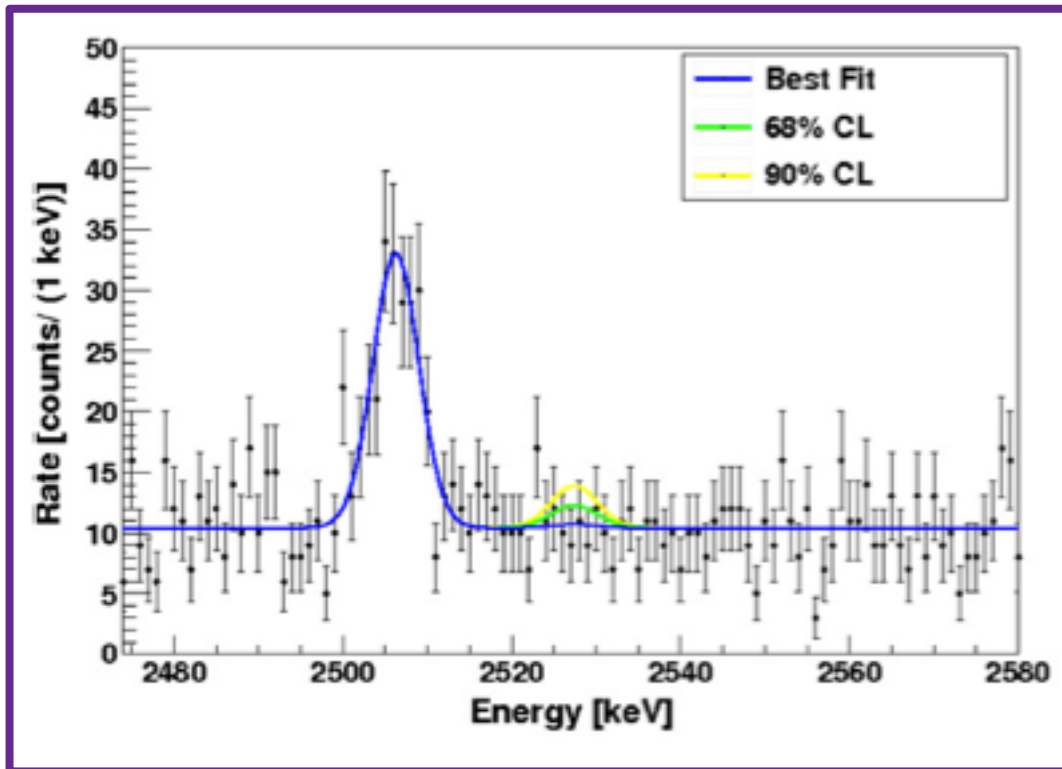
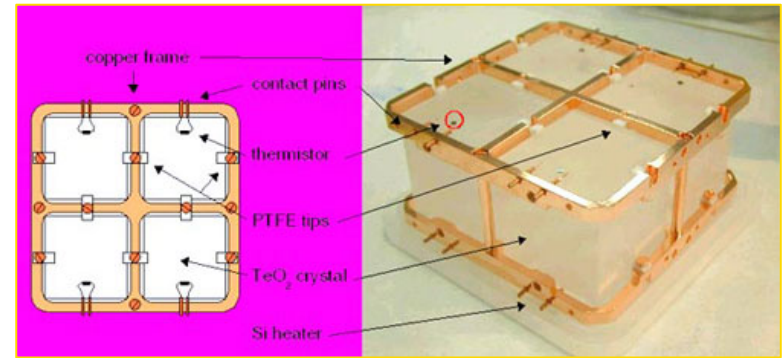
$$T_{1/2} > 1.1 \times 10^{24} \text{ years}$$

See S Blot's talk

CUORICINO

- ^{130}Te Bolometer

Astropart.Phys. 34, 822 (2011)



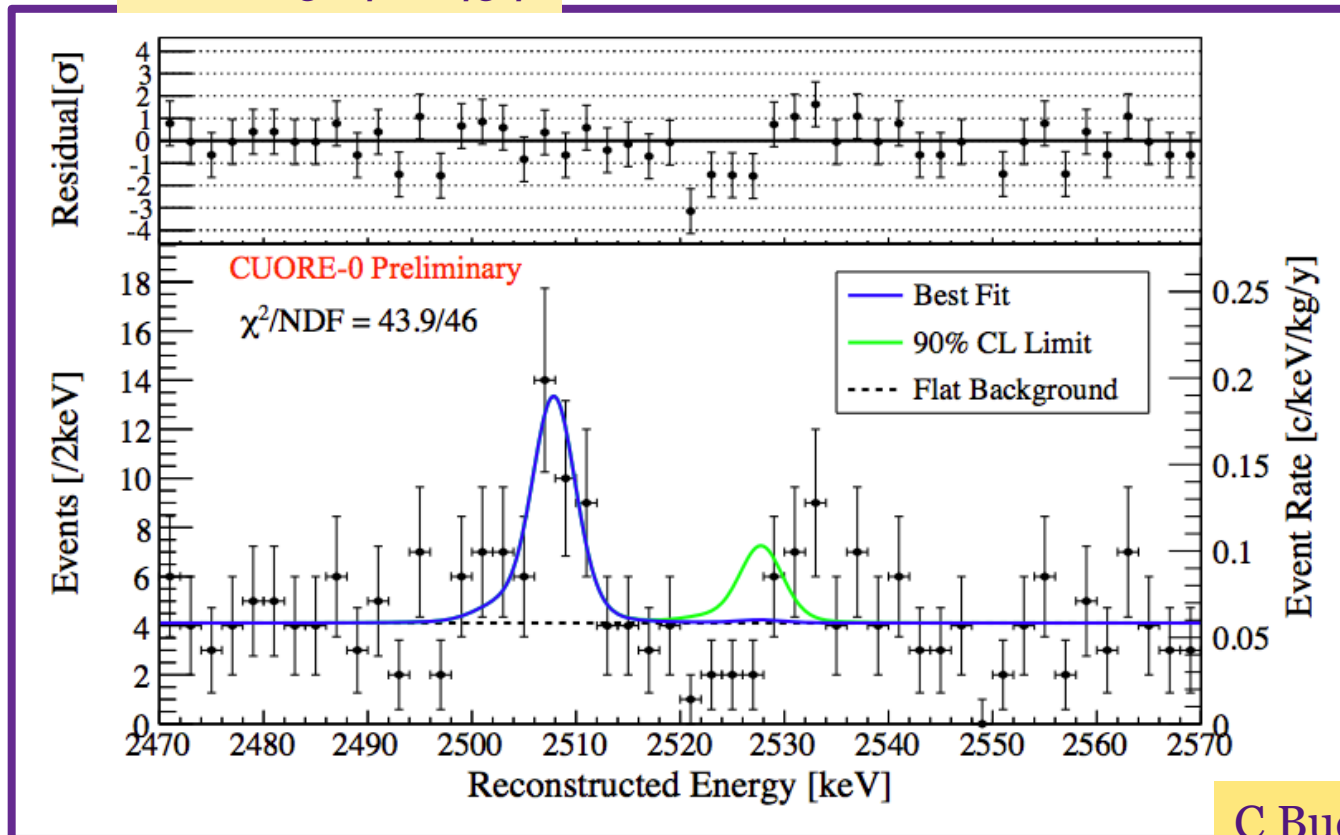
$$T_{1/2} > 2.8 \times 10^{24} \text{ years}$$

CUORE-0

- Next phase of CUORICINO technique

$$T_{1/2} > 2.7 \times 10^{24} \text{ years}$$

arXiv:1504.02454

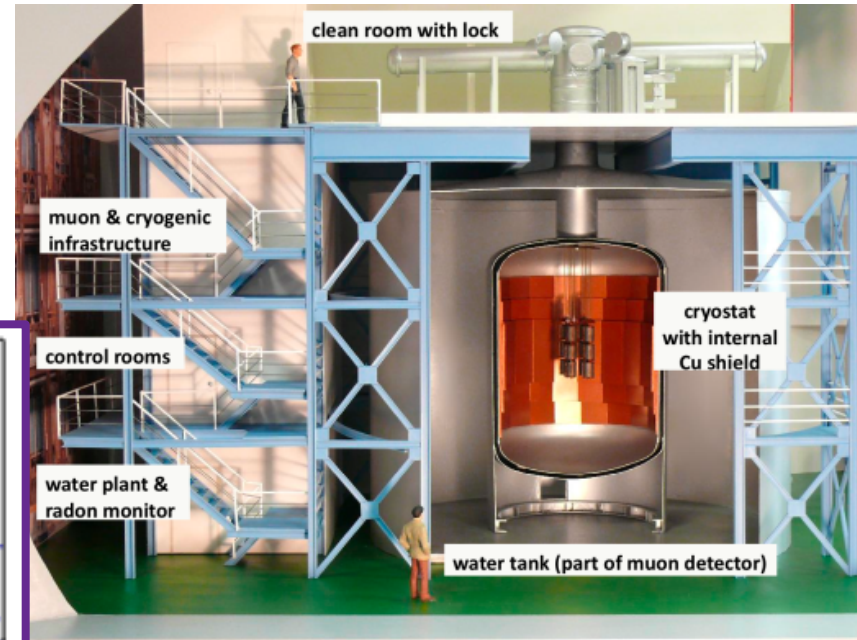
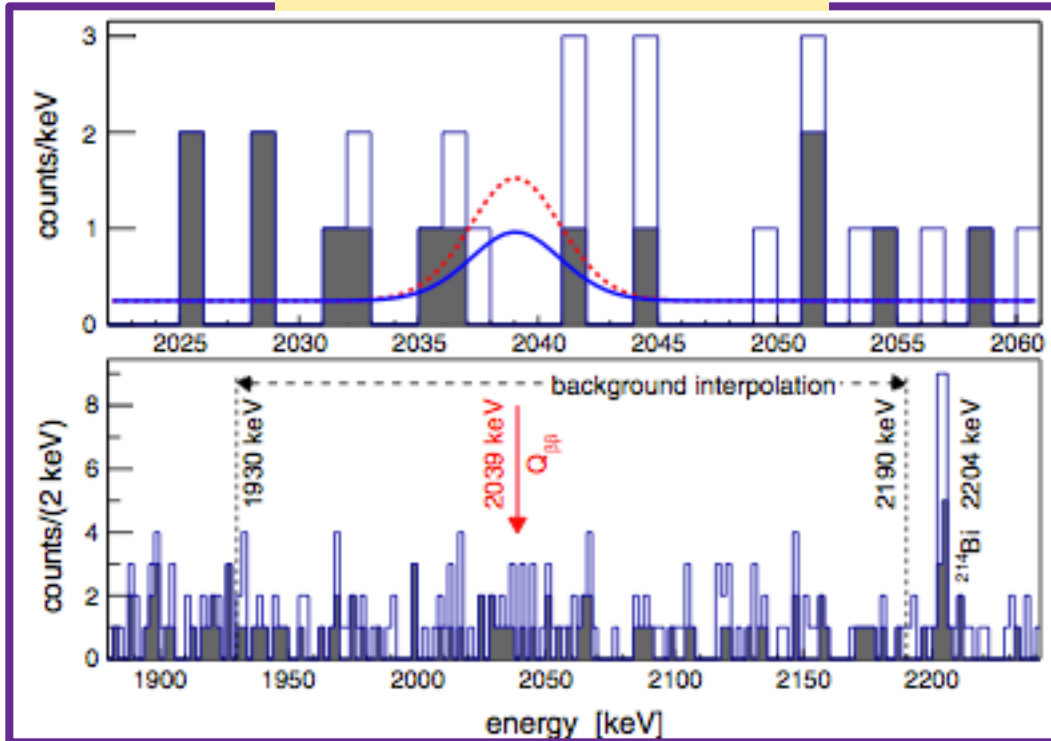


C Bucci's talk

GERDA

- Germanium crystal
 - ^{76}Ge isotope

PRL 111, 122503 (2013)



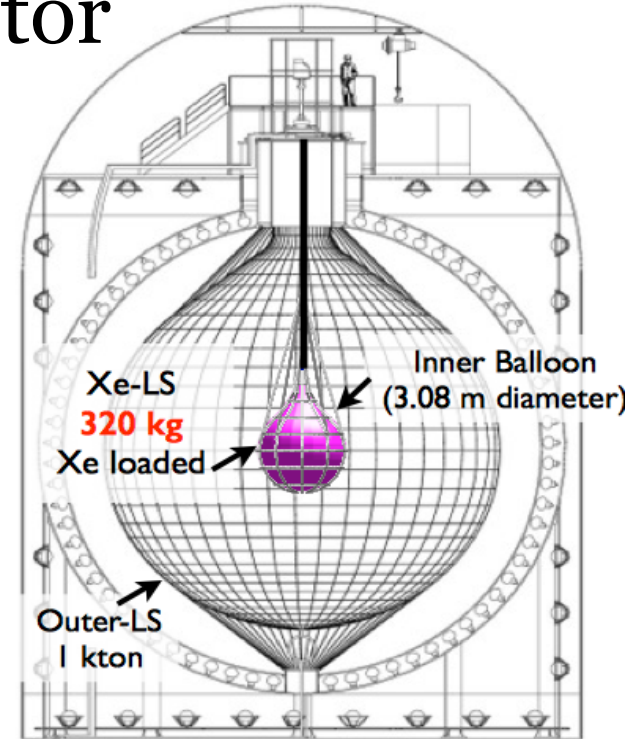
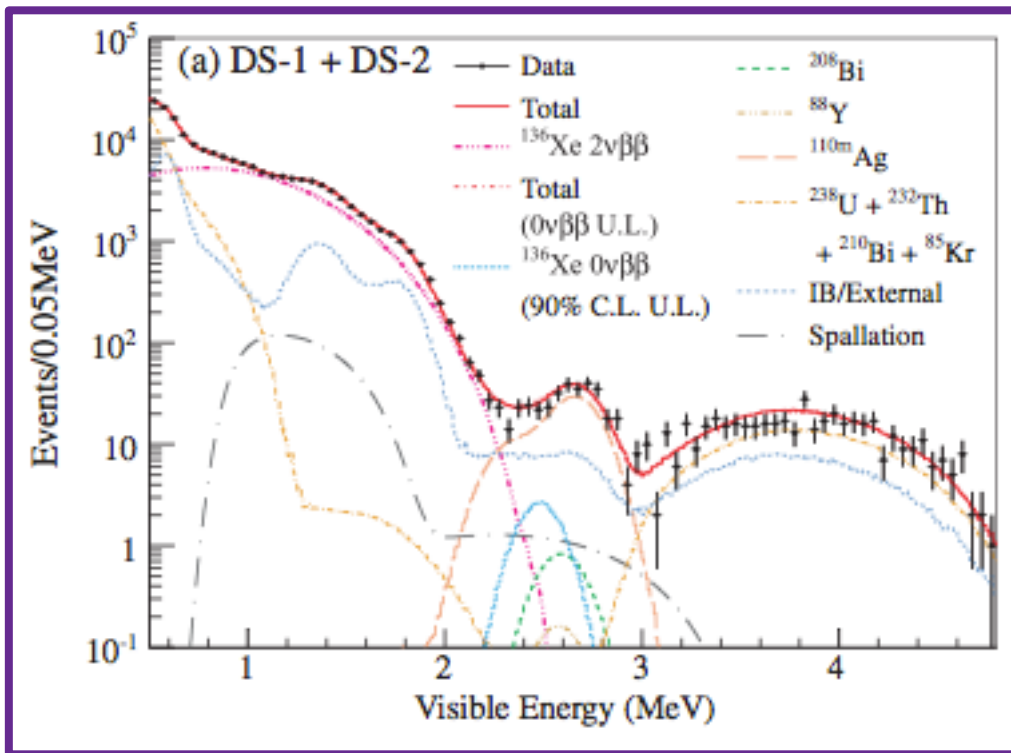
$$T_{1/2} > 2.1 \times 10^{25} \text{ years}$$

Talk by K Gusev

KamLAND-Zen

- Xenon-loaded liquid scintillator
 - ^{136}Xe isotope

PRL 110 062502 (2013)

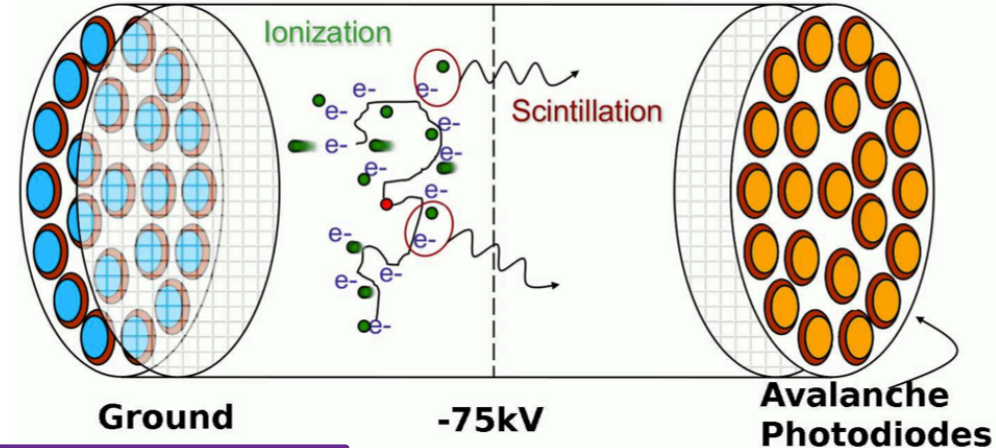


$$T_{1/2} > 1.9 \times 10^{25} \text{ years}$$

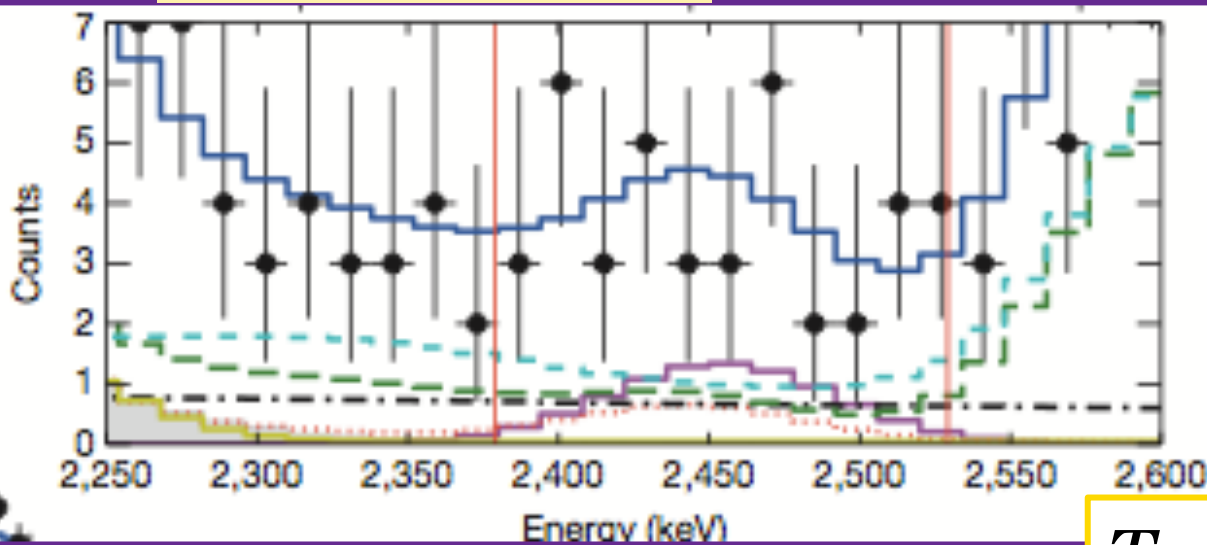
EXO-200

- Liquid xenon TPC
 - ^{136}Xe isotope

Nature 510, 229 (2014)



Talk by I Ostrovskiy



$$T_{1/2} > 1.1 \times 10^{25} \text{ years}$$

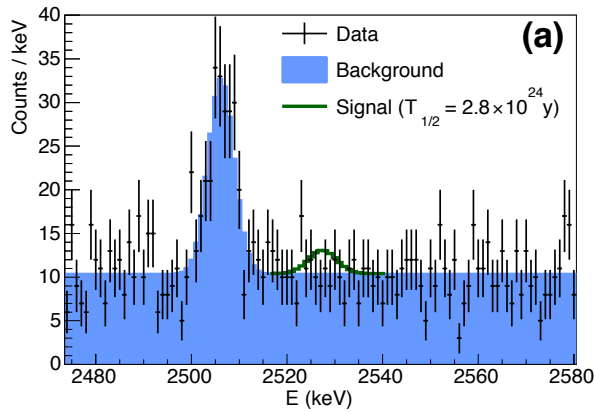
Summary of experiments

Experiment	Isotope	Method	Exposure (kg y)	$T_{1/2}$ 90% CL limit (years)	$m_{\beta\beta}$ limit (meV)
CUORICINO	^{130}Te	Bolometer	19.75	2.8×10^{24}	300 - 710
CUORE-o	^{130}Te	Bolometer	9.8	2.7×10^{24}	330 - 790
EXO-200	^{136}Xe	TPC	100	1.1×10^{25}	190 - 450
GERDA	^{76}Ge	Calorimeter	21.6	2.1×10^{25}	240 - 480
KamLAND-Zen	^{136}Xe	Loaded scintillator	89.5	1.9×10^{25}	160 - 330
NEMO-3	^{100}Mo	Foil-Tracker-Calorimeter	34.7	1.1×10^{24}	300 - 900

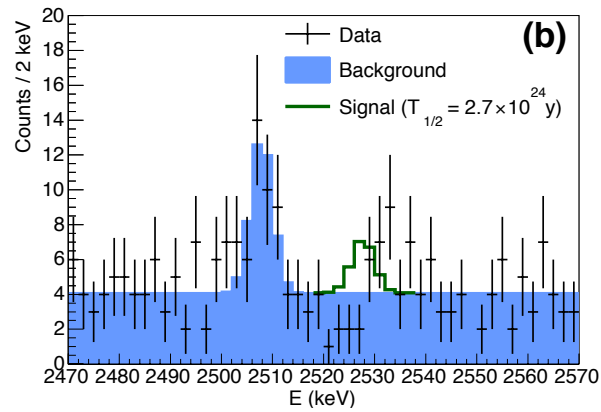
THE COMBINATION

The experimental data distributions

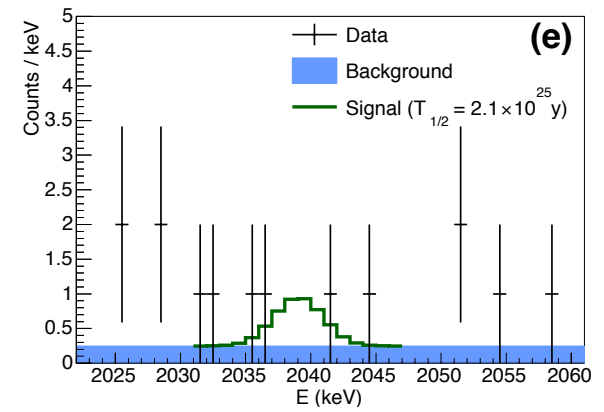
CUORICINO



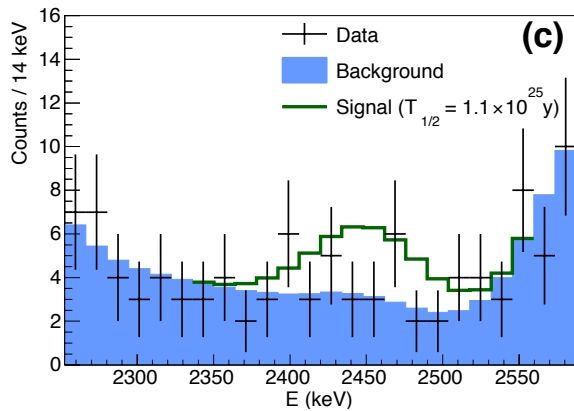
CUORE-0



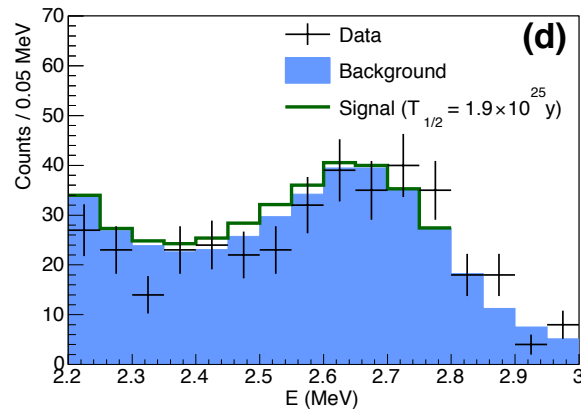
GERDA



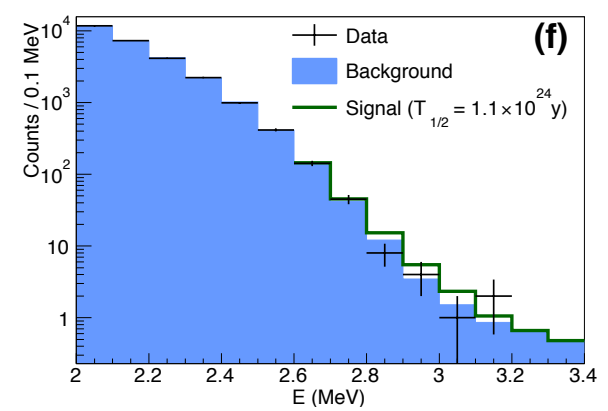
EXO



KamLAND-Zen

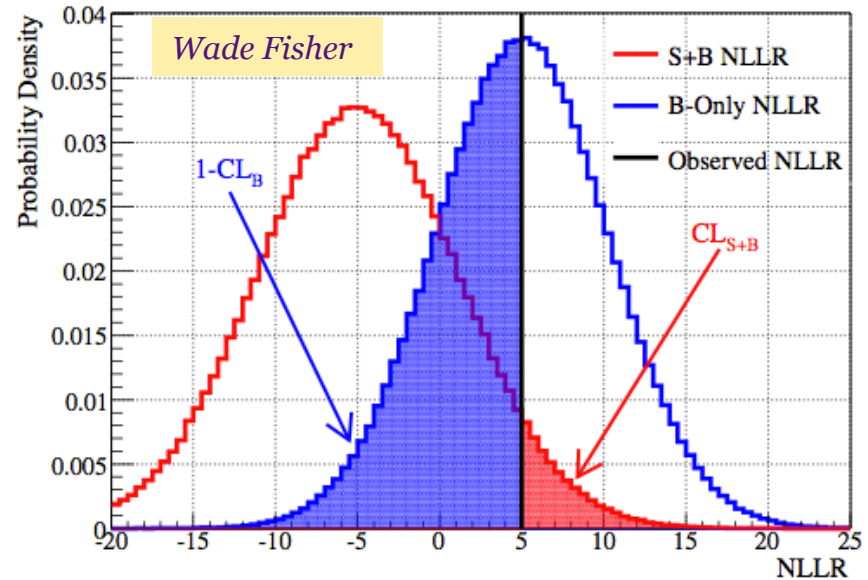


NEMO-3



Aside: CL_s limit setting method

- For each experimental data distribution there are two models
 - Background only (B)
 - Signal + background (S+B)
 - Signal normalisation can be scaled
- Calculate the Log Likelihood Ratio (LLR) for the data in both hypotheses
- Adjust signal normalisation until $CL_s = 0.1$ (for 90% CL)



$$CL_s = CL_{s+b} / CL_b$$

CL_H is the p-value of the data LLR in the H hypothesis

Cross checks of experimental inputs

CUORICINO

CUORE-o

GERDA

Experiment	Published limit (10^{24} y)	Our limit (10^{24} y)	Sensitivity (10^{24} y)	1σ range
CUORE(INO+o)	4.0	4.4	4.3	2.9 – 6.2
EXO	11	13	21	14 – 30
GERDA	21	20	21	14 – 29
KamLAND-Zen	19	17	11	7 – 15
NEMO-3 (^{100}Mo)	1.1	1.1	0.9	0.6 – 1.4

For the experiments based on ^{76}Ge and ^{100}Mo , **our limits** are in excellent agreement with those **published** by the experiment collaborations

Up to 15% difference in the ^{130}Te and ^{136}Xe experiments

Calculating half life

Decay rate :-
signal
normalisation

$$[T_{1/2}^{0\nu}]^{-1}$$

Phase space factor
(isotope-dependent)

$$= G^{0\nu} |M^{0\nu}|^2$$

$$\frac{m_{\beta\beta}^2}{m_e^2}$$

effective neutrino mass

Nuclear matrix element
(isotope-dependent)

Nuclear Matrix Elements

- Use phase space factors and NMEs to give you the relative normalisation of signals between experiments

PHYSICAL REVIEW D 79, 053001 (2009)

Quasiparticle random phase approximation uncertainties and their correlations in the analysis of $0\nu\beta\beta$ decay

PHYSICAL REVIEW C 85, 034316 (2012)

PRL 105, 252503 (2010)

PHYSICAL REVIEW LETTERS

week ending
17 DECEMBER 2010

Energy Density Functional Study of Nuclear Matrix Elements for Neutrinoless $\beta\beta$ Decay

Nuclear Physics A 818 (2009) 139–151

Disassembling the nuclear matrix elements of the neutrinoless $\beta\beta$ decay

PHYSICAL REVIEW C 87, 014315 (2013)

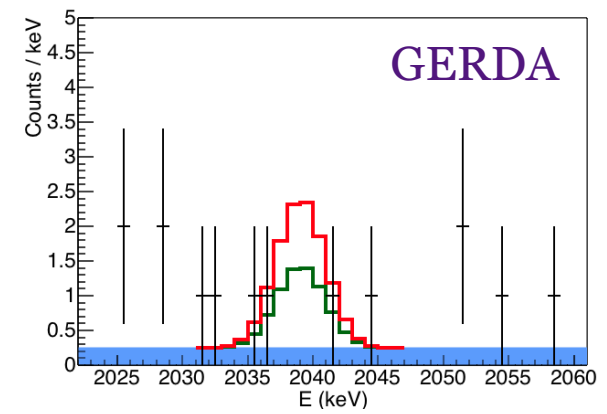
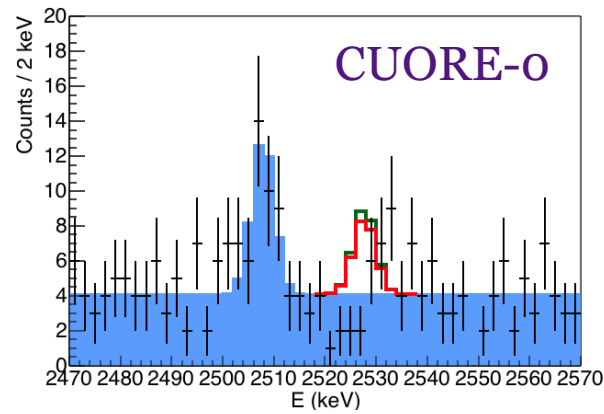
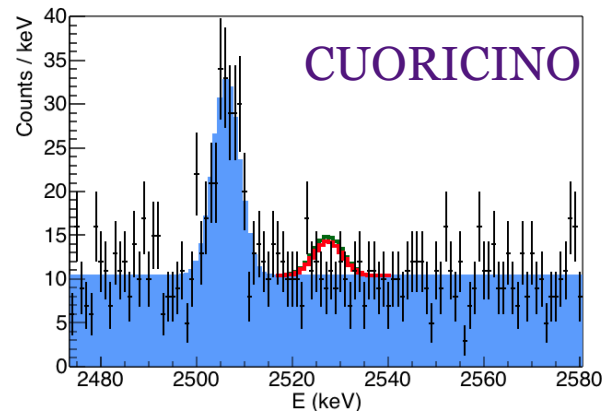
Nuclear matrix elements for double- β decay

PHYSICAL REVIEW C 87, 045501 (2013)

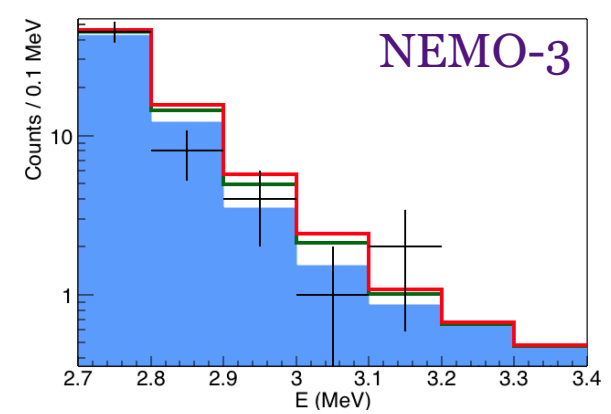
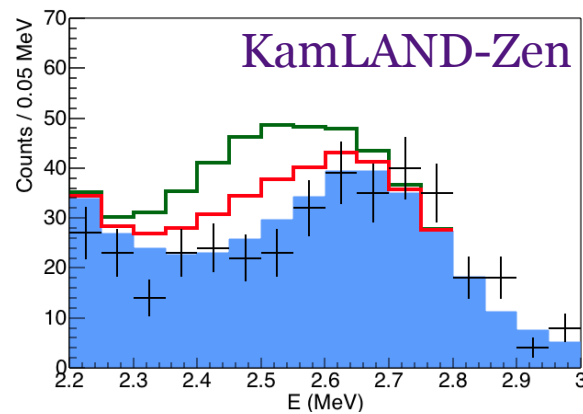
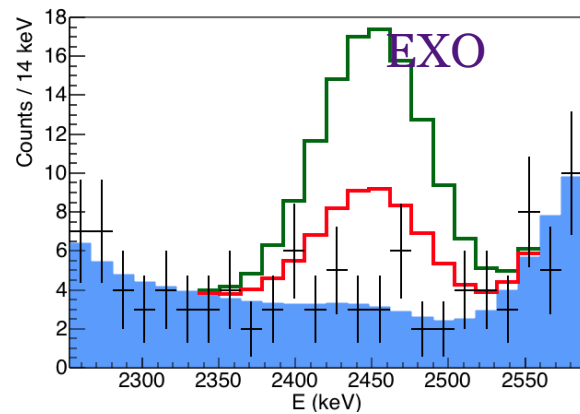
$0\nu\beta\beta$ and $2\nu\beta\beta$ nuclear matrix elements, quasiparticle random-phase approximation, and isospin symmetry restoration

Isotope	Phase Space Factor $G^{0\nu}$ (10^{-14}y^{-1})	Nuclear Matrix Element							QRPA
		RQRPA							
		GCM	IBM	NSM	A-old	A-new	B-old	B-new	
^{76}Ge	0.615	4.60	5.42	2.30	5.812	5.157	6.228	5.571	4.315
^{100}Mo	4.142	5.08	3.73	—	5.696	5.402	6.148	5.850	3.184
^{130}Te	3.699	5.13	4.03	2.12	4.306	3.888	4.810	4.373	3.148
^{136}Xe	3.793	4.20	3.33	1.76	2.437	2.177	2.735	2.460	1.795

NME effects on signal rates



For the *GCM* or *QRPA* NMEs with $m_{\beta\beta} = 400$ meV



NME effects on signal rates

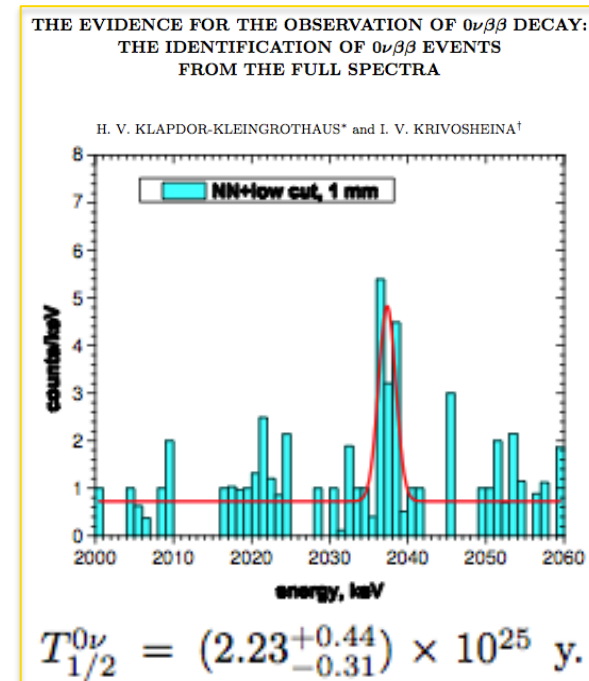
- For a constant $m_{\beta\beta}$, between two models:
 - For some isotopes (^{76}Ge , ^{100}Mo) the signal normalisation has increased
 - For other isotopes (^{130}Te , ^{136}Xe) the signal normalisation has decreased
 - Some isotopes have large change, some very similar
- We produce limits on $m_{\beta\beta}$ that depend on the NME
- The combined limit is made through summing the LLRs of the individual experiments in the CL_s method
 - signal normalisation depending on $m_{\beta\beta}$ as per the half-life formula

Results

(improvement wrt best individual experiment)

NME			limit (meV)	Improvement limit	sensitivity	HM p-value
GCM [1]			130	12%	10%	0.0001
IBM [2]			190	16%	13%	0.021
NSM [3]			310	14%	10%	0.003
QRPA [4]	Argonne	new	200	23%	25%	0.095
		old	180	26%	25%	0.100
	Bonn	new	180	28%	24%	0.073
		old	170	28%	23%	0.077
pnQRPA [5]			170	19%	16%	0.029
(R)QRPA [6]			290	23%	21%	0.311

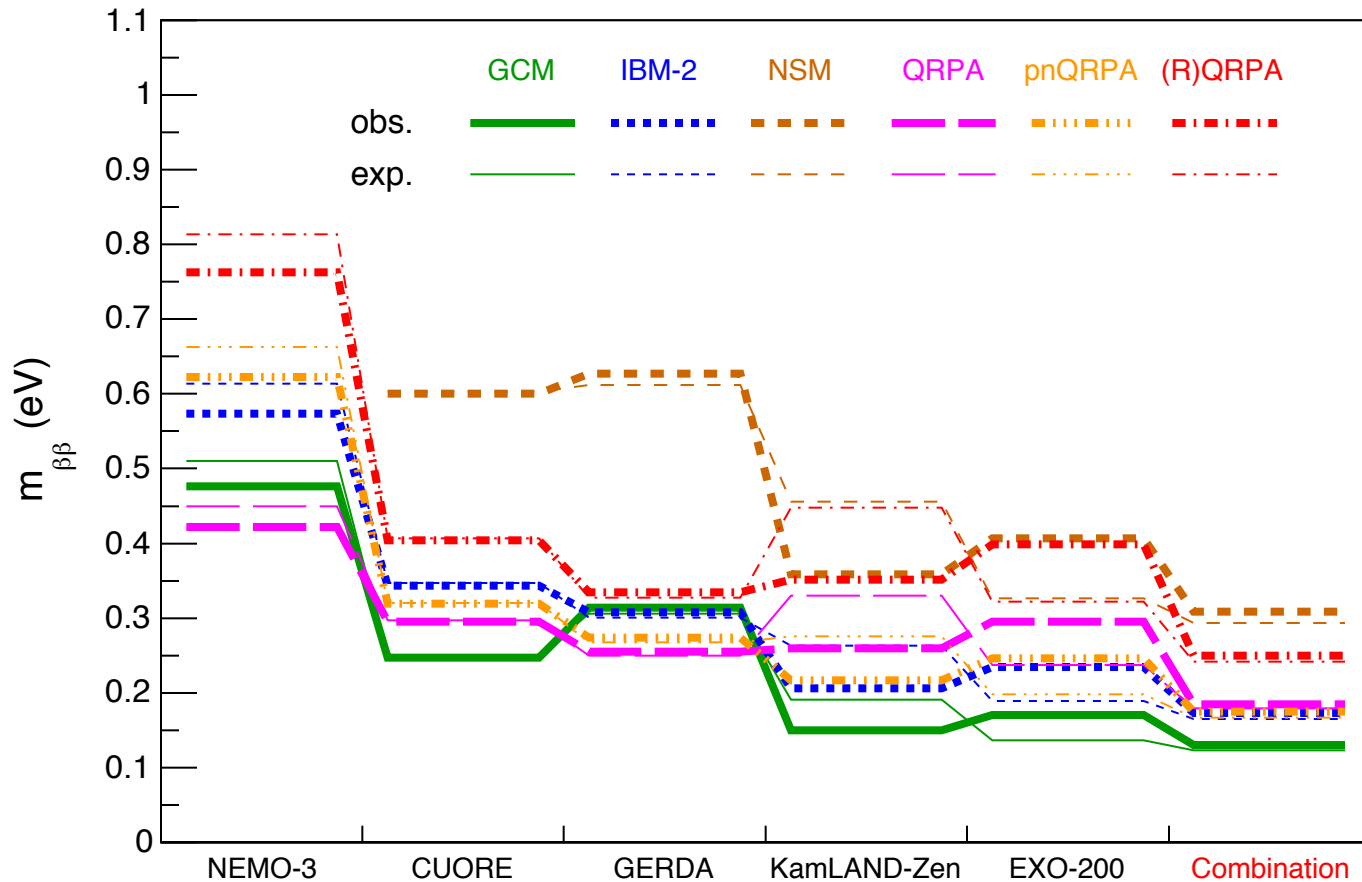
Heidelberg-Moscow experiment



Mod Phys Lett A 21 1547 (2006)

- [1] PRL 105, 252503 (2010) [2] PRC 91, 034304 (2015) [3] Nucl.Phys.A 818, 139 (2009)
 [4] PRC 87, 045501 (2013) [5] PRC 91, 024613 (2015) [6] PRD 79, 053001 (2009)

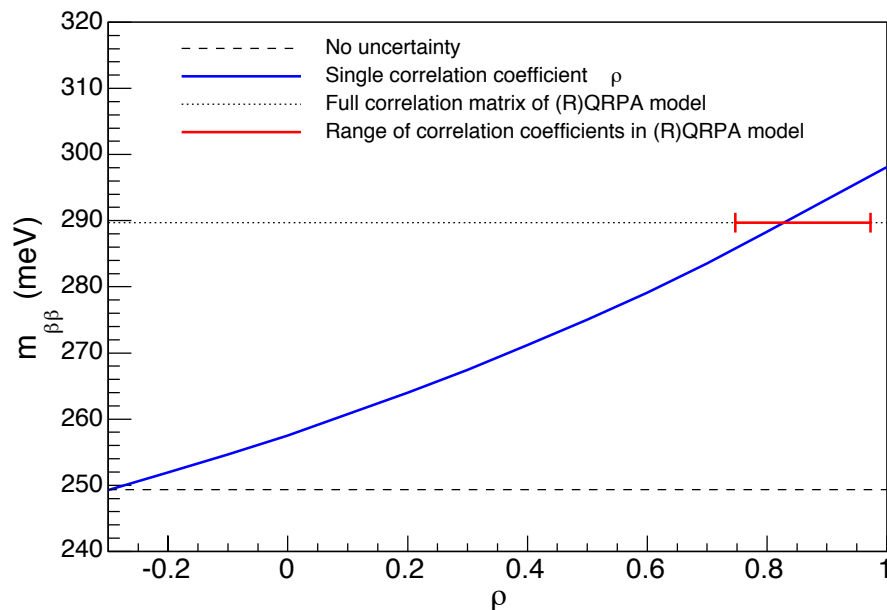
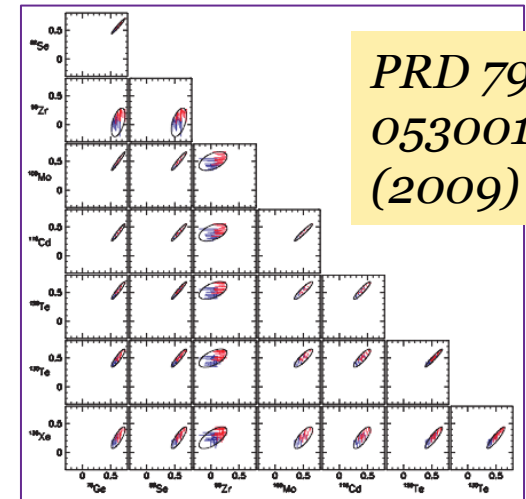
Results Graphically



QRPA uncertainty & correlation

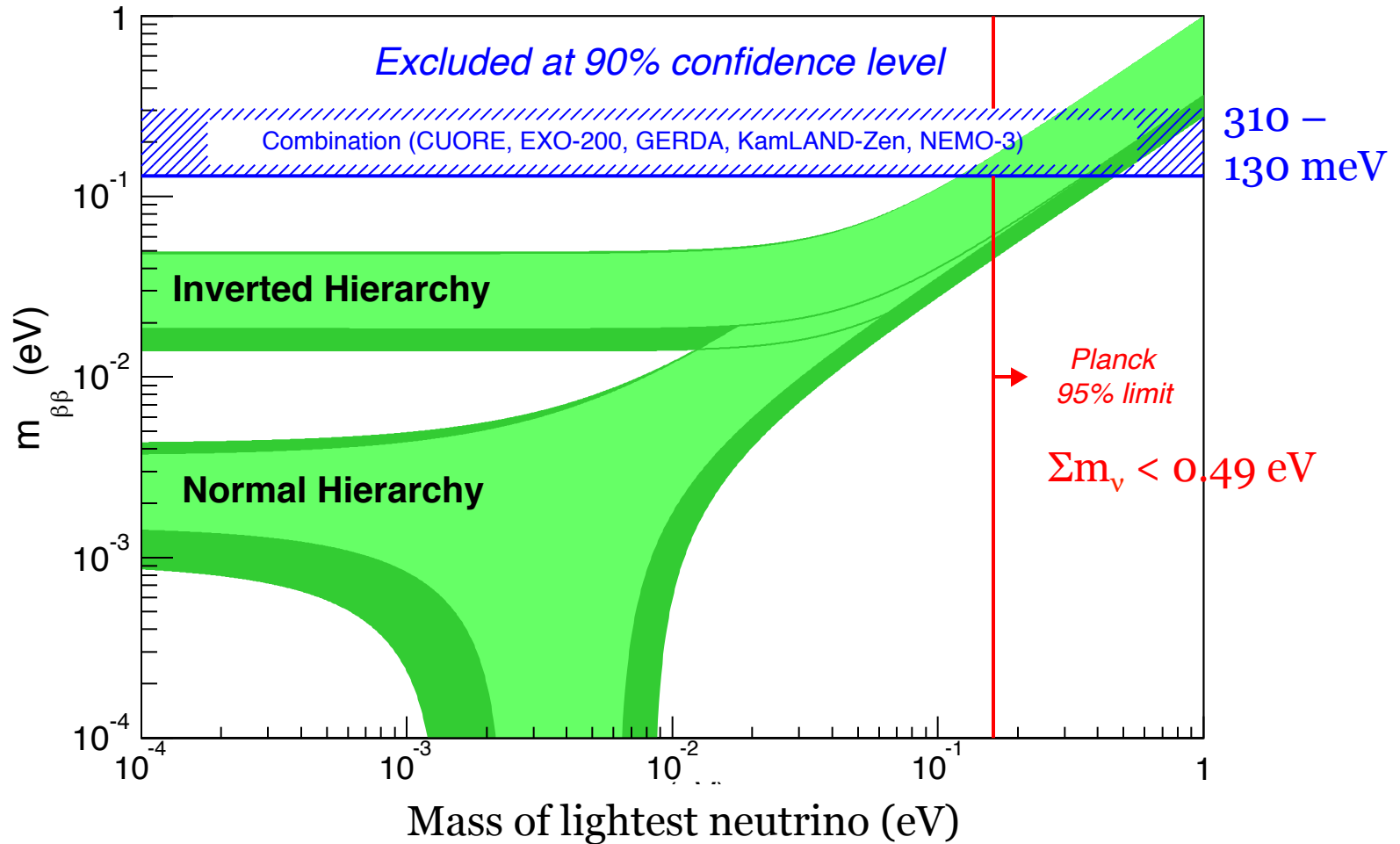
Isotope

	Rel. Unc.	Correlation Matrix			
		^{76}Ge	^{100}Mo	^{130}Te	^{136}Xe
^{76}Ge	0.191	1			
^{100}Mo	0.254	0.973	1		
^{130}Te	0.247	0.899	0.862	1	
^{136}Xe	0.293	0.805	0.747	0.916	1



- NME uncertainties lead to ~40-60% normalisation uncertainty on signal
- Effect of correlations can increase mass limit by up to 50 meV

The bigger picture



IMPLICATION FOR A STERILE MAJORANA NEUTRINO

How effective mass is calculated

$$m_{\beta\beta} = \left| m_1 |U_{e1}|^2 + m_2 |U_{e2}|^2 e^{i\alpha} + \right. \\ \left. + m_3 |U_{e3}|^2 e^{i\beta} + m_4 \sin^2 \theta_{14} e^{i\gamma} \right|$$

3+1 sterile neutrino model

Neutrino masses

PMNS matrix elements

Majorana phases

Effective mass can be thought of as the length of the sum of vectors in the complex plane

How effective mass is calculated

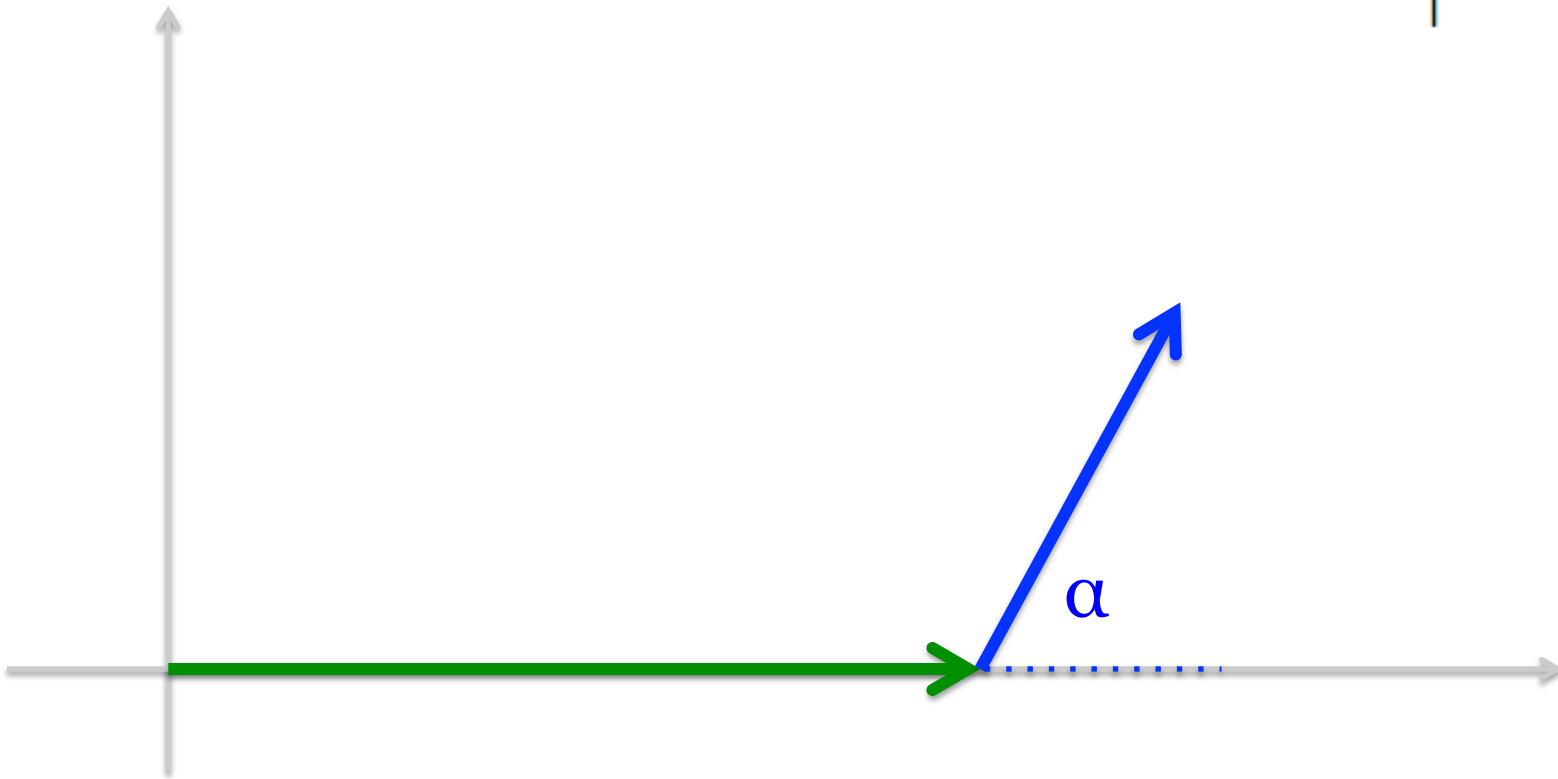
$$m_{\beta\beta} = \left| \underline{m_1 |U_{e1}|^2} + m_2 |U_{e2}|^2 e^{i\alpha} + \right. \\ \left. + m_3 |U_{e3}|^2 e^{i\beta} + m_4 \sin^2 \theta_{14} e^{i\gamma} \right|$$

Complex plane



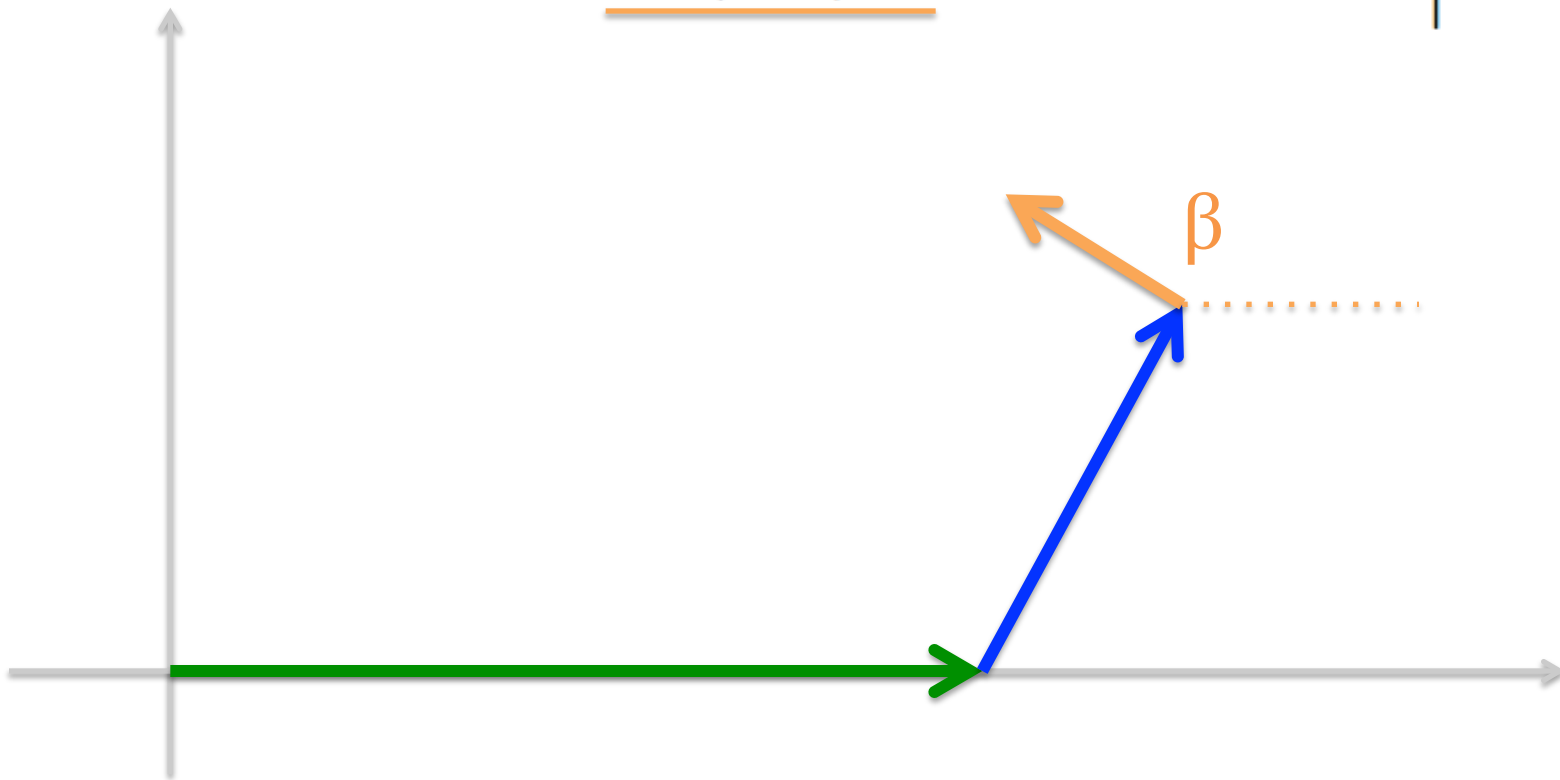
How effective mass is calculated

$$m_{\beta\beta} = \left| m_1 |U_{e1}|^2 + \underline{m_2 |U_{e2}|^2 e^{i\alpha}} + \right. \\ \left. + m_3 |U_{e3}|^2 e^{i\beta} + m_4 \sin^2 \theta_{14} e^{i\gamma} \right|$$



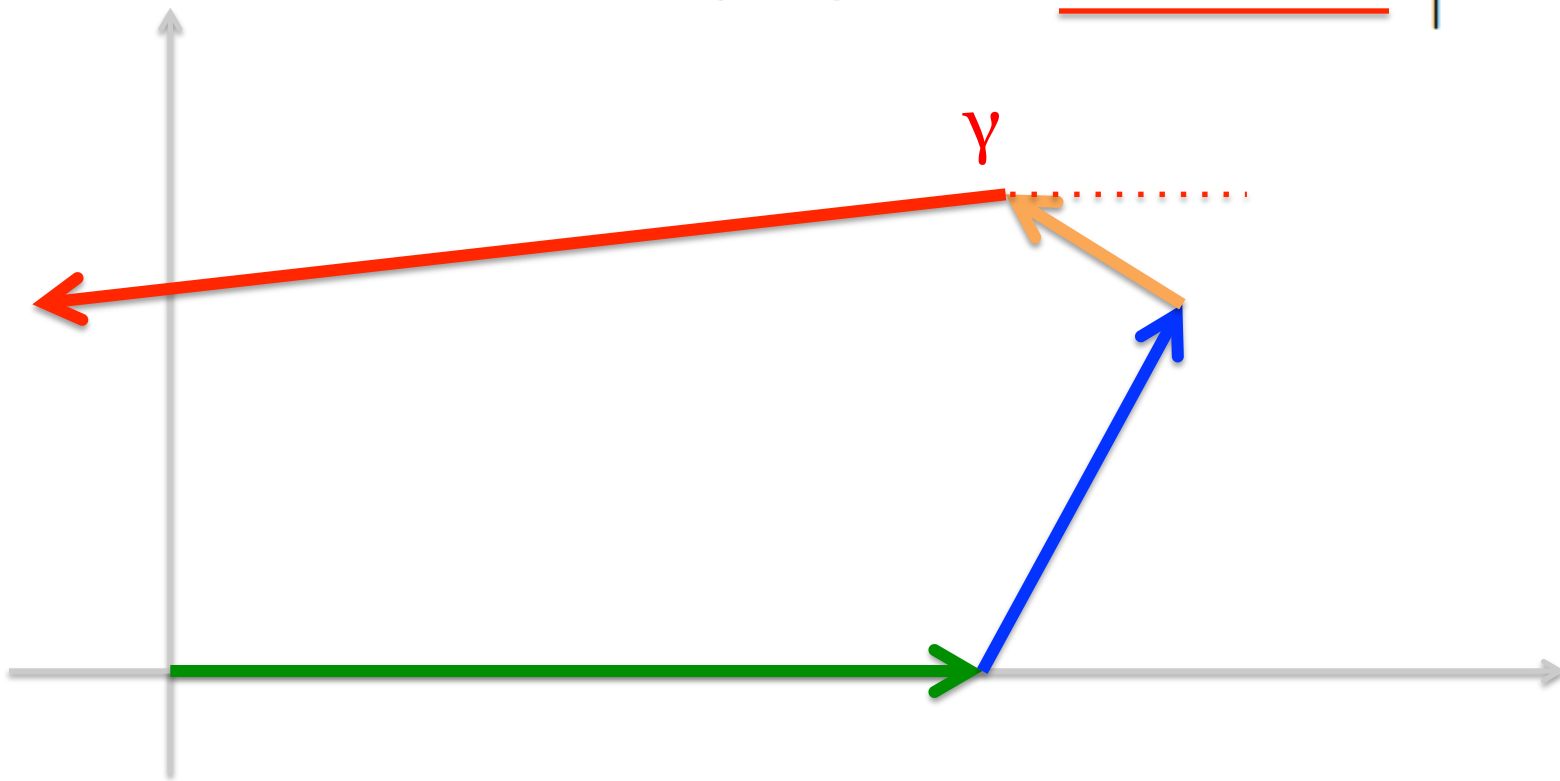
How effective mass is calculated

$$m_{\beta\beta} = \left| m_1 |U_{e1}|^2 + m_2 |U_{e2}|^2 e^{i\alpha} + \right. \\ \left. + \underline{m_3 |U_{e3}|^2 e^{i\beta}} + m_4 \sin^2 \theta_{14} e^{i\gamma} \right|$$



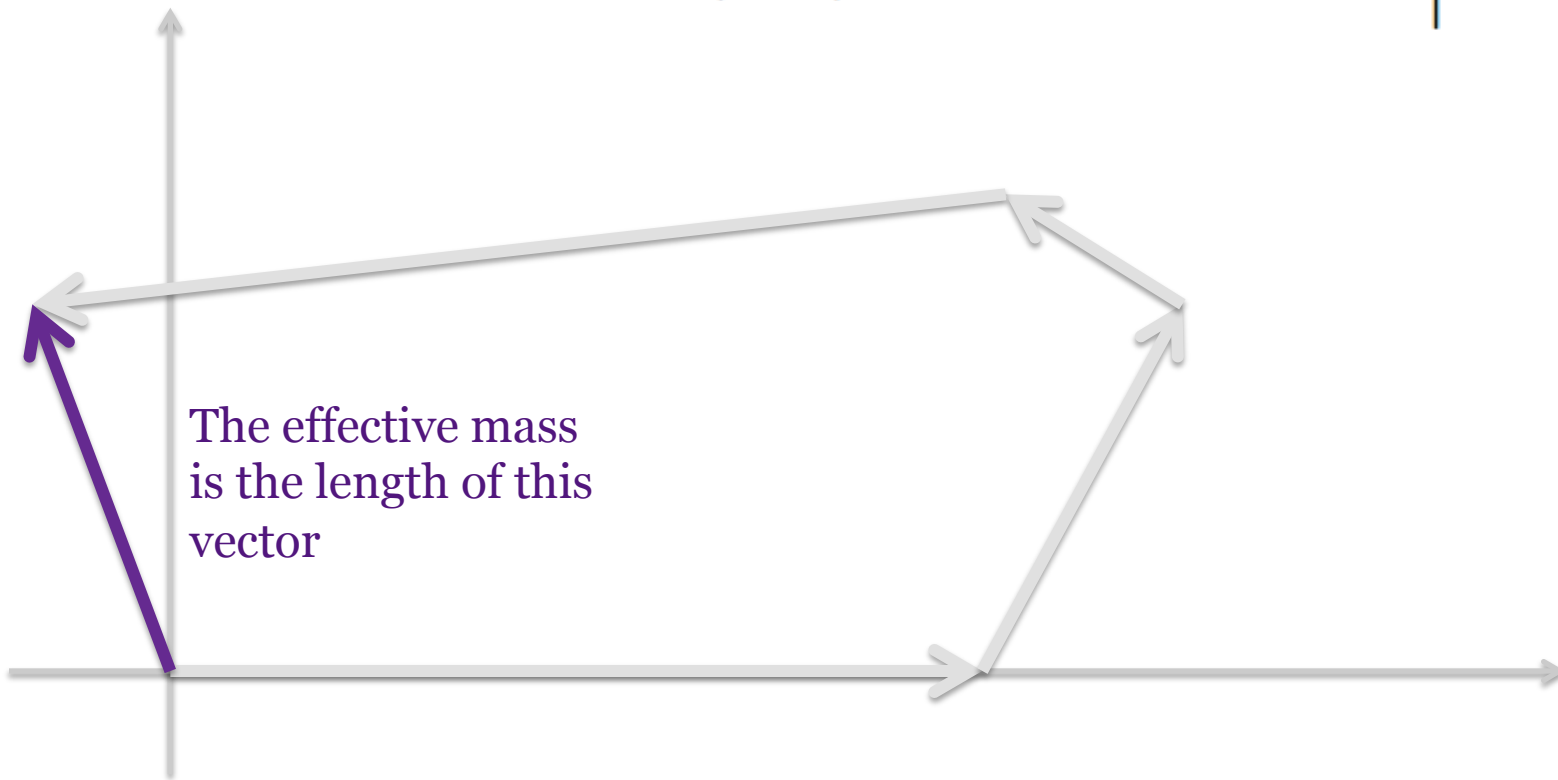
How effective mass is calculated

$$m_{\beta\beta} = \left| m_1 |U_{e1}|^2 + m_2 |U_{e2}|^2 e^{i\alpha} + \right. \\ \left. + m_3 |U_{e3}|^2 e^{i\beta} + \underline{m_4 \sin^2 \theta_{14} e^{i\gamma}} \right|$$



How effective mass is calculated

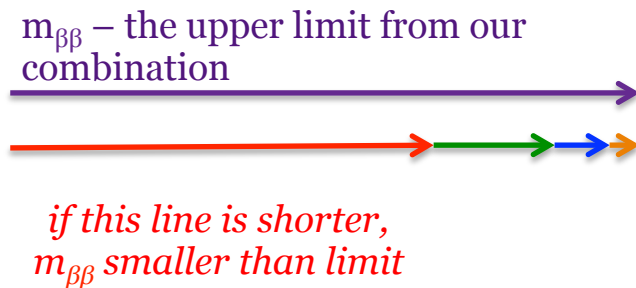
$$\underline{m_{\beta\beta}} = \left| m_1 |U_{e1}|^2 + m_2 |U_{e2}|^2 e^{i\alpha} + \right. \\ \left. + m_3 |U_{e3}|^2 e^{i\beta} + m_4 \sin^2 \theta_{14} e^{i\gamma} \right|$$



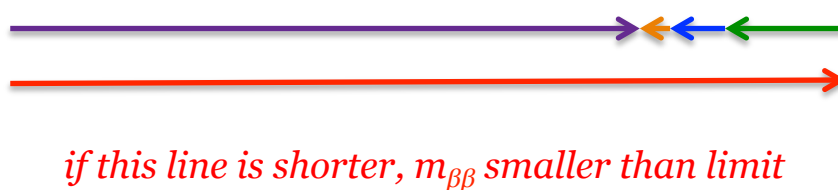
Limits on sterile neutrino

$$m_{\beta\beta} = \left| \underbrace{m_1 |U_{e1}|^2}_{\text{green}} + \underbrace{m_2 |U_{e2}|^2 e^{i\alpha}}_{\text{blue}} + \underbrace{m_3 |U_{e3}|^2 e^{i\beta}}_{\text{orange}} + \underbrace{m_4 \sin^2 \theta_{14} e^{i\gamma}}_{\text{red}} \right|$$

If we want to set limits on the **sterile neutrino**, we have two extreme cases:

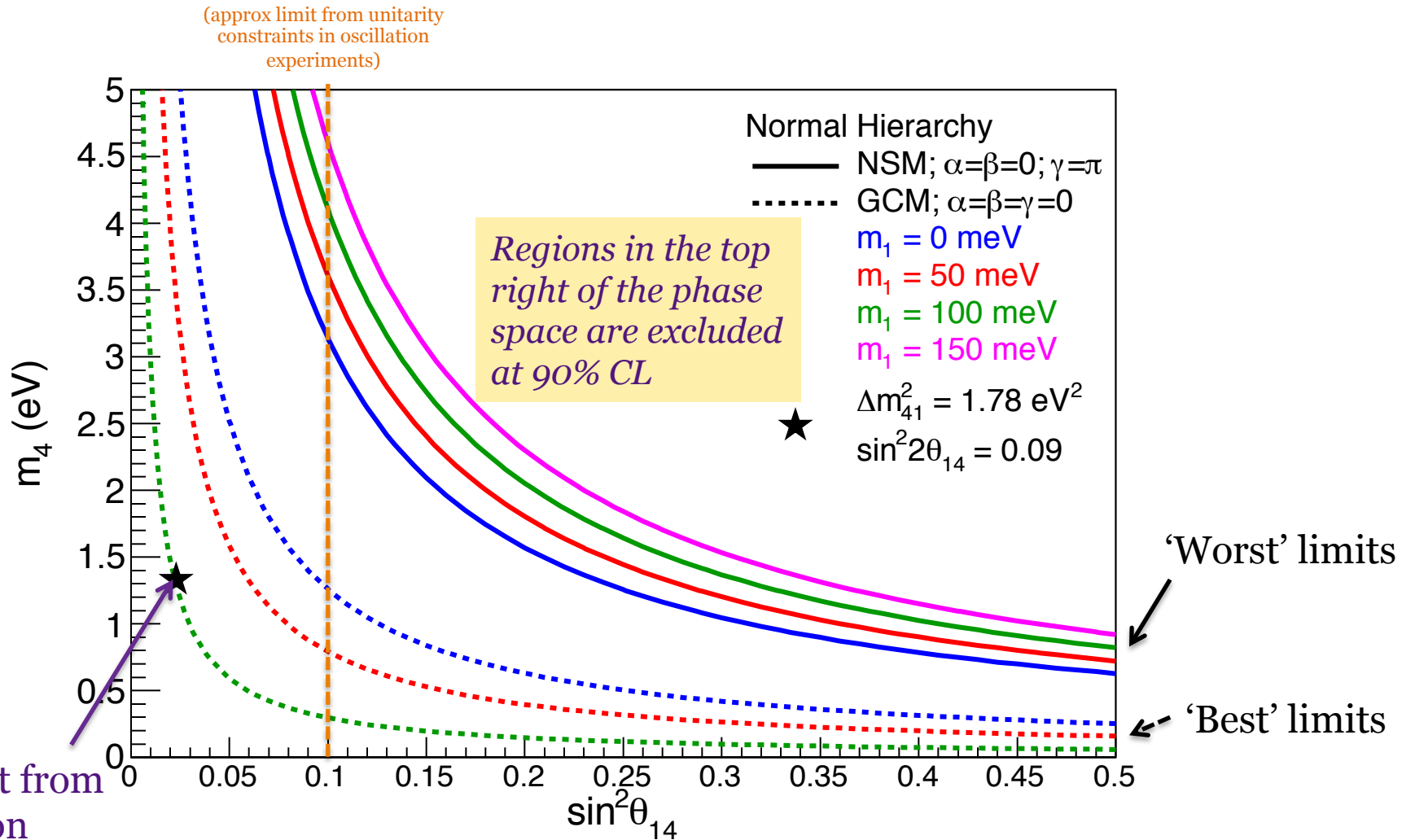


‘Best’ upper limit on m_4 for a **fixed** $m_{\beta\beta}$ – all Majorana phases aligned



‘Worst’ upper limit on m_4 – ‘sterile’ Majorana phase antiparallel to active phases

Sterile Majorana neutrino limits



Global fit from
 oscillation
 experiments
JHEP 05, 050 (2013)

SUMMARY

Summary

- A first combination of $0\nu\beta\beta$ limits with multiple isotopes has been performed
 - $m_{\beta\beta} < 130 - 310 \text{ meV}$ (depending on NME)
- The combination can offer a sensitivity improvement equivalent to $\sim 2x$ increase of exposure of the most sensitive individual experiment
- Limits can be set for a sterile Majorana neutrino, with the global best-fit candidate excluded in some favourable scenarios

- For more detailed information, please read
Phys Rev D **92, 012002**